

Strong converse theorem for state redistribution using Rényi entropies

arXiv:1506.02635

Felix Leditzky^a, Mark M. Wilde^b, Nilanjana Datta^a

a: University of Cambridge, b: Louisiana State University



UNIVERSITY OF
CAMBRIDGE



LOUISIANA STATE UNIVERSITY

ISIT Barcelona, 15 July 2016

Table of Contents

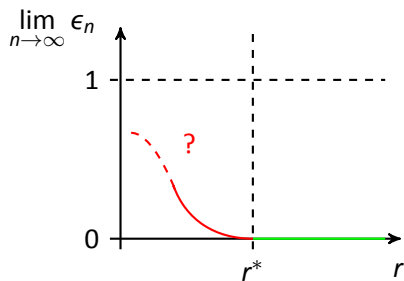
- 1 Introduction: Weak vs. strong converse
- 2 State redistribution
- 3 Rényi entropies
- 4 Main result: Strong converse for state redistribution
- 5 Summary and open questions

Table of Contents

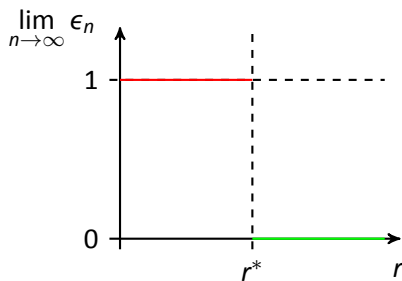
- 1 Introduction: Weak vs. strong converse
- 2 State redistribution
- 3 Rényi entropies
- 4 Main result: Strong converse for state redistribution
- 5 Summary and open questions

Introduction: Weak vs. strong converse

- ▶ Consider: information-theoretic task with optimal rate r^* .
- ▶ Code with rate r , blocklength n , and error ϵ_n .



Weak converse



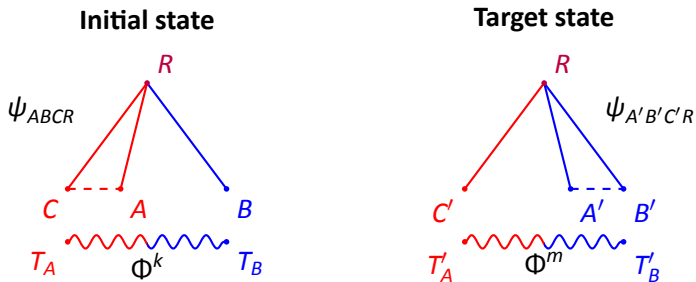
Strong converse

- ▶ This talk: Strong converse for **state redistribution**.

Table of Contents

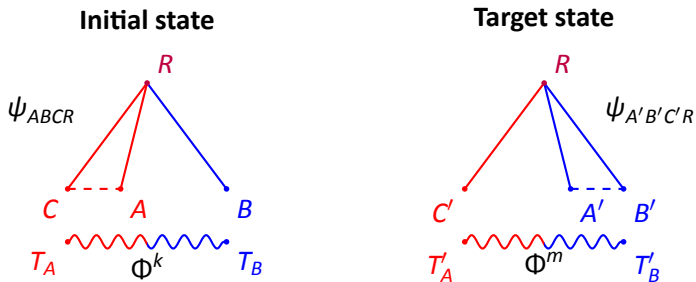
- 1 Introduction: Weak vs. strong converse
- 2 State redistribution**
- 3 Rényi entropies
- 4 Main result: Strong converse for state redistribution
- 5 Summary and open questions

State redistribution



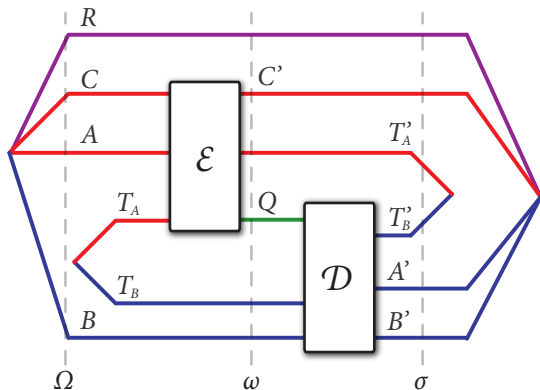
- ▶ Consider a tripartite state ρ_{ABC} shared between Alice (AC) and Bob (B), with purification ψ_{ABCR} where R is a reference system.
- ▶ Furthermore, Alice and Bob share a maximally entangled state $\Phi_{T_A T_B}^k$ of Schmidt rank k .

State redistribution



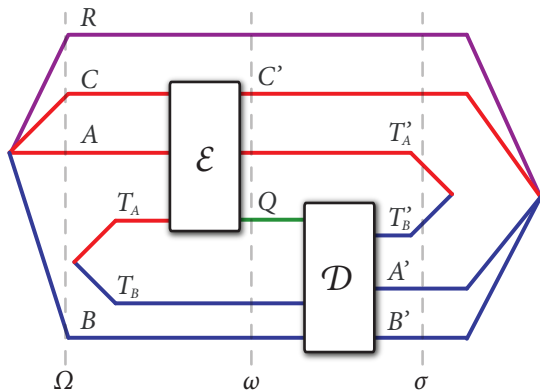
- ▶ **State redistribution:** Transfer A -part of ρ from Alice to Bob.
- ▶ Use shared entanglement, local encoding \mathcal{E} (Alice) and decoding \mathcal{D} (Bob), and quantum communication Q from Alice to Bob.

State redistribution



- ▶ Use shared entanglement, local encoding \mathcal{E} (Alice) and decoding \mathcal{D} (Bob), and quantum communication Q from Alice to Bob.

State redistribution



- ▶ State redistribution **generalizes many quantum information-processing tasks** like Schumacher compression, coherent state merging, and quantum state splitting.

State redistribution

- ▶ Consider now **n i.i.d. copies** of initial state ρ , and a shared maximally entangled state Φ^{k_n} of Schmidt rank k_n .
- ▶ **Target state:** n i.i.d. copies of ρ , and some maximally entangled state Φ^{m_n} of Schmidt rank m_n .
- ▶ State redistribution protocol outputs state σ_n .
- ▶ **Figure of merit:** Fidelity

$$F_n := F(\sigma_n, \Phi^{m_n} \otimes \psi^{\otimes n}),$$

where $F(\omega, \tau) := \|\omega^{1/2}\tau^{1/2}\|_1$.

State redistribution

► **Operational quantities:**

▷ Quantum communication cost $q_n := \frac{1}{n} \log |Q^n|$

▷ Entanglement cost

$$e_n := \frac{1}{n}(\log k_n - \log m_n) = \frac{1}{n}(\log |\mathcal{T}_A^n| - \log |\mathcal{T}'_A^n|)$$

► Note that e_n can be negative, in which case entanglement is gained in the protocol rather than consumed.

Definition (Achievable rates)

A rate pair (e, q) is achievable if there exists a sequence of state redistribution protocols such that

$$\limsup_{n \rightarrow \infty} e_n = e, \quad \limsup_{n \rightarrow \infty} q_n = q, \quad \lim_{n \rightarrow \infty} F_n = 1.$$

State redistribution: Optimal rates

Theorem (Optimal rates)

Luo and Devetak (2009) and Yard and Devetak (2009)

The rate pair (e, q) is achievable if and only if

$$q + e \geq S(A|B)_\rho \quad \text{and} \quad q \geq \frac{1}{2}I(A; C|B)_\rho,$$

where $S(A|B)_\rho = S(AB)_\rho - S(B)_\rho$ is the conditional entropy and $I(A; C|B)_\rho = S(A|B)_\rho - S(A|BC)_\rho$ is the conditional mutual information.

State redistribution: Optimal rates

- ▶ Preceding theorem proves a **weak converse**:

$$\limsup(q_n + e_n) < S(A|B)_\rho \quad \vee \quad \limsup q_n < \frac{1}{2}I(A; C|B)_\rho \\ \implies \lim F_n < 1$$

- ▶ Our goal is to prove **strong converse**, that is,

$$F_n \leq \exp(-Kn)$$

for some suitable constant $K > 0$, yielding $\lim F_n = 0$.

- ▶ Note: A strong converse theorem has also been obtained by [Berta et al. \(2016\)](#) using smooth entropies.
- ▶ We use the so-called Rényi entropy method (cf. [Arimoto \(1973\)](#) and [Ogawa and Nagaoka \(1999\)](#)).

Table of Contents

- 1 Introduction: Weak vs. strong converse
- 2 State redistribution
- 3 Rényi entropies**
- 4 Main result: Strong converse for state redistribution
- 5 Summary and open questions

Rényi entropies

Definition (Sandwiched Rényi divergence)

Müller-Lennert et al. (2013) and Wilde et al. (2014)

Let ρ be a quantum state and σ be positive semidefinite with $\text{supp } \rho \subseteq \text{supp } \sigma$. For $\alpha \in (0, 1) \cup (1, \infty)$, we set $\gamma = (1 - \alpha)/2\alpha$ and define the α -sandwiched Rényi divergence as

$$\tilde{D}_\alpha(\rho \parallel \sigma) := \frac{1}{\alpha - 1} \log \text{Tr} (\sigma^\gamma \rho \sigma^\gamma)^\alpha.$$

Quantum generalization of the classical Rényi divergence

$$D_\alpha(P \parallel Q) = \frac{1}{\alpha - 1} \log \sum_x P(x)^\alpha Q(x)^{1-\alpha}.$$

That is, $\tilde{D}_\alpha(\rho_X \parallel \sigma_X) = D_\alpha(P \parallel Q)$ for $\rho_X = \bigoplus_x P(x)$, $\sigma_X = \bigoplus_x Q(x)$.

Rényi entropies

Properties of the sandwiched Rényi divergence:

- ▶ Recovers the quantum relative entropy:

$$\lim_{\alpha \rightarrow 1} \tilde{D}_\alpha(\rho \parallel \sigma) = D(\rho \parallel \sigma),$$

where $D(\rho \parallel \sigma) = \text{Tr}(\rho(\log \rho - \log \sigma))$.

- ▶ **Additivity:**

$$\tilde{D}_\alpha(\rho_1 \otimes \rho_2 \parallel \sigma_1 \otimes \sigma_2) = \tilde{D}_\alpha(\rho_1 \parallel \sigma_1) + \tilde{D}_\alpha(\rho_2 \parallel \sigma_2)$$

- ▶ **Data processing inequality:** Frank and Lieb (2013)

For every quantum operation Λ and $\alpha \geq 1/2$,

$$\tilde{D}_\alpha(\rho \parallel \sigma) \geq \tilde{D}_\alpha(\Lambda(\rho) \parallel \Lambda(\sigma)).$$

Rényi entropies

Definition (Derived entropies)

Let ρ_{AB} be a bipartite quantum state with marginal ρ_A . Then we define the following quantities:

▶ **Rényi entropy:** $S_\alpha(A)_\rho := -\tilde{D}_\alpha(\rho_A \| \mathbb{1}_A)$;

▶ **Rényi conditional entropy:**

$$\tilde{S}_\alpha(A|B)_\rho := -\min_{\sigma_B} \tilde{D}_\alpha(\rho_{AB} \| \mathbb{1}_A \otimes \sigma_B)$$

▶ **Rényi mutual information:**

$$\tilde{I}_\alpha(A; B)_\rho := \min_{\sigma_B} \tilde{D}_\alpha(\rho_{AB} \| \rho_A \otimes \sigma_B).$$

These recover the von Neumann quantities in the limit $\alpha \rightarrow 1$, and satisfy additivity and data processing inequality as well.

Table of Contents

- 1 Introduction: Weak vs. strong converse
- 2 State redistribution
- 3 Rényi entropies
- 4 Main result: Strong converse for state redistribution**
- 5 Summary and open questions

Main result: Strong converse theorem

Theorem

Let $\alpha \in (1/2, 1)$ and β be such that $1/\alpha + 1/\beta = 2$. Then for all $n \in \mathbb{N}$ we have the following bounds on $F_n = F(\sigma_n, \Phi^{m_n} \otimes \psi^{\otimes n})$:

$$F_n \leq \exp \left\{ -n\kappa(\alpha) [S_\beta(AB)_\rho - S_\alpha(B)_\rho - (q_n + e_n)] \right\}$$

$$F_n \leq \exp \left\{ -n\kappa(\alpha) \left[\tilde{S}_\beta(R|B)_\rho - \tilde{S}_\alpha(R|AB)_\rho - 2q_n \right] \right\}$$

$$F_n \leq \exp \left\{ -n\kappa(\alpha) \left[\tilde{I}_\alpha(R; AB)_\rho - \tilde{I}_\beta(R; B)_\rho - 2q_n \right] \right\}$$

where $\kappa(\alpha) = (1 - \alpha)/2\alpha$, and q_n and e_n are the quantum communication cost and entanglement cost, respectively.

Main result: Strong converse theorem

$$F_n \leq \exp \left\{ -n\kappa(\alpha) \left[S_{\beta}(AB)_{\rho} - S_{\alpha}(B)_{\rho} - (q_n + e_n) \right] \right\}$$

$$F_n \leq \exp \left\{ -n\kappa(\alpha) \left[\tilde{S}_{\beta}(R|B)_{\rho} - \tilde{S}_{\alpha}(R|AB)_{\rho} - 2q_n \right] \right\}$$

$$F_n \leq \exp \left\{ -n\kappa(\alpha) \left[\tilde{I}_{\alpha}(R; AB)_{\rho} - \tilde{I}_{\beta}(R; B)_{\rho} - 2q_n \right] \right\}$$

- ▶ Recall that state redistribution is possible iff $q + e \geq S(A|B)_{\rho}$.
- ▶ We have $S_{\beta}(AB)_{\rho} - S_{\alpha}(B)_{\rho} \xrightarrow{\alpha \rightarrow 1} S(A|B)_{\rho}$.
- ▶ Hence, if $q_n + e_n < S(A|B)_{\rho}$, there is $\alpha_0 < 1$ such that

$$K(\alpha_0) := \kappa(\alpha_0) \left[S_{\beta(\alpha_0)}(AB)_{\rho} - S_{\alpha_0}(B)_{\rho} - (q_n + e_n) \right] > 0,$$

and $F_n \leq \exp(-nK(\alpha_0))$.

Main result: Strong converse theorem

$$F_n \leq \exp \left\{ -n\kappa(\alpha) [S_\beta(AB)_\rho - S_\alpha(B)_\rho - (q_n + e_n)] \right\}$$

$$F_n \leq \exp \left\{ -n\kappa(\alpha) \left[\tilde{S}_\beta(R|B)_\rho - \tilde{S}_\alpha(R|AB)_\rho - 2q_n \right] \right\}$$

$$F_n \leq \exp \left\{ -n\kappa(\alpha) \left[\tilde{I}_\alpha(R; AB)_\rho - \tilde{I}_\beta(R; B)_\rho - 2q_n \right] \right\}$$

- ▶ Recall that state redistribution is possible iff $2q \geq I(A; C|B)_\rho$.
- ▶ $\tilde{S}_\beta(R|B)_\rho - \tilde{S}_\alpha(R|AB)_\rho \xrightarrow{\alpha \rightarrow 1} I(A; C|B)_\rho$ and
 $\tilde{I}_\alpha(R; AB)_\rho - \tilde{I}_\beta(R; B)_\rho \xrightarrow{\alpha \rightarrow 1} I(A; C|B)_\rho$.
- ▶ Hence, if $2q < I(A; C|B)_\rho$, there is $\alpha_0 < 1$ and $K(\alpha_0) > 0$ such that $F_n \leq \exp(-nK(\alpha_0))$.

Key ingredient in the proof of the main result

Theorem (Fidelity bound)

Let ρ_{AB} and σ_{AB} be bipartite quantum states, and for $\alpha \in (1/2, 1)$ let β be such that $1/\alpha + 1/\beta = 2$, then

$$(i) \quad S_\alpha(A)_\rho - S_\beta(A)_\sigma \geq \kappa(\alpha)^{-1} \log F(\rho_A, \sigma_A)$$

$$(ii) \quad \tilde{S}_\alpha(A|B)_\rho - \tilde{S}_\beta(A|B)_\sigma \geq \kappa(\alpha)^{-1} \log F(\rho_{AB}, \sigma_{AB}).$$

If $\rho_A = \sigma_A$, then also:

$$(iii) \quad \tilde{I}_\beta(A; B)_\rho - \tilde{I}_\alpha(A; B)_\sigma \geq \kappa(\alpha)^{-1} \log F(\rho_{AB}, \sigma_{AB}).$$

Here, $\kappa(\alpha) = (1 - \alpha)/2\alpha$ as before.

- ▶ Generalizes van Dam and Hayden (2002), who proved (i).
- ▶ Our proof method: Hölder inequality for Schatten p -norms.

Table of Contents

- 1 Introduction: Weak vs. strong converse
- 2 State redistribution
- 3 Rényi entropies
- 4 Main result: Strong converse for state redistribution
- 5 Summary and open questions**

Summary

- ▶ We established a full strong converse theorem for state redistribution.
- ▶ Employed “Rényi entropy method” based on entropic quantities derived from sandwiched Rényi divergence, together with fidelity bounds.
- ▶ Same method yields strong converse for:
 - ▷ measurement compression with quantum side information
 - ▷ randomness extraction
 - ▷ data compression with quantum side information

(see full paper at [arXiv:1506.02635](https://arxiv.org/abs/1506.02635))

Open questions

- ▶ Are these bounds optimal?
→ Rényi quantities as strong converse exponents?
- ▶ Can we use the fidelity bounds for proving strong converse theorems for other tasks?

Thank you very much!

Arimoto, S. (1973). **IEEE Trans. Inf. Th.** 19.3, pp. 357–359.

Berta, M. et al. (2016). **IEEE Trans. Inf. Th.** 62.3, pp. 1425–1439. arXiv: 1409.4338
[quant-ph].

Frank, R. L. and E. H. Lieb (2013). **Journal of Mathematical Physics** 54.12, p. 122201. arXiv:
1306.5358 [math-ph].

Luo, Z. and I. Devetak (2009). **IEEE Trans. Inf. Th.** 55.3, pp. 1331–1342. arXiv:
quant-ph/0611008.

Müller-Lennert, M. et al. (2013). **Journal of Mathematical Physics** 54.12, p. 122203. arXiv:
1306.3142 [quant-ph].

Ogawa, T. and H. Nagaoka (1999). **IEEE Trans. Inf. Th.** 45.7, pp. 2486–2489. arXiv:
quant-ph/9808063.

Wilde, M. M. et al. (2014). **Comm. Math. Phys.** 331.2, pp. 593–622. arXiv: 1306.1586
[quant-ph].

Yard, J. T. and I. Devetak (2009). **IEEE Trans. Inf. Th.** 55.11, pp. 5339–5351. arXiv: 0706.2907
[quant-ph].

van Dam, W. and P. Hayden (2002). **arXiv preprint**. arXiv: quant-ph/0204093.