## ASTR 3740 Relativity & Cosmology Spring 2025. Problem Set 4. Due Wed 5 Mar

Warning: this problem set is quite lengthy, so please do not wait until the last day to start

### 1. Trajectories of particles in the Schwarzschild geometry

In this problem you will find it helpful to visit John Walker's web site at

http://www.fourmilab.to/gravitation/orbits/.

The most fun part of the site is the Java applet, so you will probably want to seek out a Java-enabled computer, although you can also use John Walker's site without Java.

In what follows, the time t, radial coordinate r, polar angle  $\theta$ , and azimuthal angle  $\phi$  are the usual Schwarzschild coordinates in the Schwarzschild metric (with c=1 as usual)

$$ds^{2} = -\left(1 - \frac{r_{s}}{r}\right)dt^{2} + \frac{dr^{2}}{\left(1 - \frac{r_{s}}{r}\right)} + r^{2}\left(d\theta^{2} + \sin^{2}\theta \,d\phi^{2}\right) , \qquad (1.1)$$

with  $r_s$  the Schwarzschild radius

$$r_s = 2GM (1.2)$$

Without loss of generality, the trajectory of a particle falling freely in the Schwarzschild geometry may be taken to lie in the equatorial plane,  $\theta = \pi/2$ . For a particle of finite (nonzero) mass, the trajectory satisfies the equations

$$\left(1 - \frac{r_s}{r}\right) \frac{dt}{d\tau} = E \,\,, \tag{1.3a}$$

$$r^2 \frac{d\phi}{d\tau} = L , \qquad (1.3b)$$

$$r^{2} \frac{d\phi}{d\tau} = L , \qquad (1.3b)$$

$$\left(\frac{dr}{d\tau}\right)^{2} + V_{\text{eff}} = E^{2} , \qquad (1.3c)$$

where  $\tau$  is the proper time of the particle, and E and L are constants, the particle's energy and angular momentum per unit mass. The quantity  $V_{\text{eff}}$  is the effective potential given by

$$V_{\text{eff}} = \left(1 - \frac{r_s}{r}\right) \left(1 + \frac{L^2}{r^2}\right) . \tag{1.4}$$

#### (a) (3 points) Check

Are John Walker's equations the same as the ones given above (aside from possible differences in notation)?

### (b) (3 points) Velocity at infinity

Argue from equations (1.3) that relative to the rest frame of the Schwarzschild geometry, the radial velocity  $v_r \equiv dr/dt$  and the transverse velocity  $v_{\perp} \equiv r d\phi/dt$  (the  $\equiv$  sign means "is defined to be equal to") of the particle at extremely large distances from the Schwarzschild geometry,  $r \to \infty$ , are related to E and L by

$$v_r^2 = 1 - \frac{1}{E^2} - \frac{L^2}{E^2 r^2} \,, (1.5a)$$

$$v_{\perp} = \frac{L}{Er} \,\,, \tag{1.5b}$$

(note that L can be extremely large at large r, so L/r is not necessarily zero in the limit  $r \to \infty$ ). Hence show that the velocity  $v_{\infty} \equiv (v_r^2 + v_{\perp}^2)^{1/2}$  of the particle as  $r \to \infty$  is related to its energy E by

$$E = \frac{1}{(1 - v_{\infty}^2)^{1/2}} \ . \tag{1.6}$$

What does it mean if E < 1?

## (c) (3 points) Extrema of the effective potential

Find the radii at which the effective potential  $V_{\text{eff}}$  is a maximum or a minimum, i.e.  $dV_{\text{eff}}/dr = 0$ , as a function of angular momentum L. You should find that extrema exist only if the absolute value |L| of the angular momentum exceeds a certain critical value  $L_c$ . What is that critical value?

## (d) (3 points) Sketch

Sketch what the effective potential looks like for values of L (i) less than, (ii) equal to, (iii) greater than the critical value  $L_c$ . Make sure to label the axes clearly and correctly. Describe physically, in words, what the possible orbital trajectories are for the various cases. [Hint: You will need to experiment with different choices of axes to make the graph look good. I found it clearer *not* to start the effective potential at zero. For cases (i) and (iii), values near the critical value  $L_c$  showed the distinction most clearly.]

# (e) (3 points) Circular orbits

Circular orbits, satisfying  $dr/d\tau=0$ , occur where the effective potential is a minimum (stable orbit) or a maximum (unstable orbit). Show (from your equation for the extrema of the effective potential) that the angular momentum L of a particle in circular orbit at radius r satisfies

$$|L| = \frac{r}{\left(\frac{2r}{r_s} - 3\right)^{1/2}} \,, \tag{1.7}$$

and hence show also that the energy E in this circular orbit is

$$E = \frac{2^{1/2}(r - r_s)}{[r(2r - 3r_s)]^{1/2}} . {1.8}$$

## (f) (3 points) Orbital period

Show that the orbital period t, as measured by an observer at rest at infinity, of a particle in circular orbit at radius r is given by Kepler's 3rd law (yes, it's true even in the fully general relativistic case!)

$$\frac{GMt^2}{(2\pi)^2} = r^3 \ . {1.9}$$

[Hint: The time measured by an observer at rest at infinity is just the Schwarzschild time t. Argue that the azimuthal angle  $\phi$  evolves according to

$$\frac{d\phi}{dt} = \frac{L(r - r_s)}{Er^3} \ . \tag{1.10}$$

The period t is the time taken for  $\phi$  to change by  $2\pi$ .

### (g) (3 points) Infall time

Calculate the proper time  $\tau$  for a particle with L=0 and E=1 to fall from a finite radius r to the singularity at zero radius. What is the physical significance of the choice L=0 and E=1? [Hint: Write down the equation for  $dr/d\tau$  for L=0 and E=1, and then solve it.]

## (h) (3 points) Infall time — numbers

Use your answer to part (g) to show that the proper time to fall from the Schwarzschild radius  $r = r_s$  to the singularity (for L = 0 and E = 1) is, in units including c,

$$\tau = \frac{4GM}{3c^3} \ . \tag{1.11}$$

Evaluate your answer, in seconds, for the case of a black hole of mass  $4\times10^6~\rm M_{\odot}$ , such as may be in the center of our Galaxy, the Milky Way. [Constants: https://physics.nist.gov/cuu/Constants/  $c=299,792,458~\rm m\,s^{-1}$ ;  $G=6.6743\times10^{-11}~\rm m^3\,kg^{-1}\,s^{-2}$ ;  $1~\rm M_{\odot}=1.99\times10^{30}\,kg$ .]

### 2. Photons in the Schwarzschild geometry

The orbit equations (1.3) would appear to break down for photons, which have zero mass, hence infinite energy per unit mass E (cf. equation [1.6] for  $v_{\infty} = 1$ ) and infinite angular momentum per unit mass L. Another way of looking at this is that photons follow null geodesics,  $d\tau = 0$ , so that  $\tau$ , which does not change, is not a very useful time coordinate for expressing the equations of motion of photons.

The difficulty is cured by introducing an "affine parameter"  $\lambda = E\tau$ , which functions as a good scalar coordinate along null geodesics. In terms of the affine parameter  $\lambda$ , the equations of motion (1.3) for freely falling massless particles, such as photons, become

$$\left(1 - \frac{r_s}{r}\right) \frac{dt}{d\lambda} = 1 ,$$
(2.1a)

$$r^2 \frac{d\phi}{d\lambda} = J , \qquad (2.1b)$$

$$r^{2} \frac{d\phi}{d\lambda} = J , \qquad (2.1b)$$

$$\left(\frac{dr}{d\lambda}\right)^{2} + \mathcal{V}_{\text{eff}} = 1 , \qquad (2.1c)$$

where J = L/E is the photon's angular momentum per unit energy, and  $\mathcal{V}_{\text{eff}} = V_{\text{eff}}/E$  is the effective potential given by

$$\mathcal{V}_{\text{eff}} = \left(1 - \frac{r_s}{r}\right) \frac{J^2}{r^2} \ . \tag{2.2}$$

## (a) (6 points) Circular orbits

Circular orbits, occur where the effective potential  $\mathcal{V}_{\text{eff}}$  is a minimum (stable orbit) or a maximum (unstable orbit). At what radius can photons orbit in circles? Is the orbit stable or unstable?