

Separating Matrix-valued Quantum Correlation via Embezzlement and Teleportation

Li Gao

University of Illinois at Urbana-Champaign

Rocky Mountain Summit on Quantum Information, June 28, 2018

Joint work with Samuel J. Harris and Marius Junge

Quantum correlations

Quantum correlations are probabilistic correlations arose from quantum measurement. Suppose both Alice and Bob has d measurements (POVM) and each POVM has m values, i.e on some Hilbert space H

$$\sum_{a=1}^m E_a^x = \sum_{b=1}^m F_b^y = 1, \quad x, y = 1, \dots, d.$$

We consider the correlation matrix $[p(a, b|x, y)]_{\substack{1 \leq x, y \leq d \\ 1 \leq a, b \leq m}}$ in $(\mathbb{R}_+)^{d^2 m^2}$ given

$$p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$$

where $|\psi\rangle \in H$ is a unit vector and $[E_a^x, F_b^y] = 0$.

Quantum correlations corresponds to strategies in two player non-local games.

Quantum correlations

Quantum correlations are probabilistic correlations arose from quantum measurement. Suppose both Alice and Bob has d measurements (POVM) and each POVM has m values, i.e on some Hilbert space H

$$\sum_{a=1}^m E_a^x = \sum_{b=1}^m F_b^y = 1, \quad x, y = 1, \dots, d.$$

We consider the correlation matrix $[p(a, b|x, y)]_{\substack{1 \leq x, y \leq d \\ 1 \leq a, b \leq m}}$ in $(\mathbb{R}_+)^{d^2 m^2}$ given

$$p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$$

where $|\psi\rangle \in H$ is a unit vector and $[E_a^x, F_b^y] = 0$.

Quantum correlations corresponds to strategies in two player non-local games.

Quantum correlations

Quantum correlations are probabilistic correlations arose from quantum measurement. Suppose both Alice and Bob has d measurements (POVM) and each POVM has m values, i.e on some Hilbert space H

$$\sum_{a=1}^m E_a^x = \sum_{b=1}^m F_b^y = 1, \quad x, y = 1, \dots, d.$$

We consider the correlation matrix $[p(a, b|x, y)]_{\substack{1 \leq x, y \leq d \\ 1 \leq a, b \leq m}}$ in $(\mathbb{R}_+)^{d^2 m^2}$ given

$$p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$$

where $|\psi\rangle \in H$ is a unit vector and $[E_a^x, F_b^y] = 0$.

Quantum correlations corresponds to strategies in two player non-local games.

Quantum correlations

Quantum correlations are probabilistic correlations arose from quantum measurement. Suppose both Alice and Bob has d measurements (POVM) and each POVM has m values, i.e on some Hilbert space H

$$\sum_{a=1}^m E_a^x = \sum_{b=1}^m F_b^y = 1, \quad x, y = 1, \dots, d.$$

We consider the correlation matrix $[p(a, b|x, y)]_{\substack{1 \leq x, y \leq d \\ 1 \leq a, b \leq m}}$ in $(\mathbb{R}_+)^{d^2 m^2}$ given

$$p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$$

where $|\psi\rangle \in H$ is a unit vector and $[E_a^x, F_b^y] = 0$.

Quantum correlations corresponds to strategies in two player non-local games.

Correlation under different assumptions

A correlation matrix $\{p(a, b|x, y)\}$ is called

- **local** if $p(a, b|x, y) = \sum_n \lambda_n p_n(a|x) q_n(b|y)$ ($\dim H = 1$, hidden variable)
- **quantum** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where $|\psi\rangle \in H = H_A \otimes H_B$ with $\dim H_A, \dim H_B < \infty$, $\{E_a^x\}$ PVMs on H_A for $\{F_b^y\}$ PVMs on H_B .
(finite dimensional **separable** quantum systems)
- **quantum spatial** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where we allow $\dim H_A = \dim H_B = \infty$.
(Infinite dimensional **separable** quantum systems)
- **quantum commuting** if $p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$,
where E_a^x, F_b^y PVMs on H and $[E_a^x, F_b^y] = 0$.
(Infinite dimensional and **commuting** observables)

local \subseteq quantum \subseteq quantum spatial \subseteq quantum commuting

Correlation under different assumptions

A correlation matrix $\{p(a, b|x, y)\}$ is called

- **local** if $p(a, b|x, y) = \sum_n \lambda_n p_n(a|x) q_n(b|y)$ ($\dim H = 1$, hidden variable)
- **quantum** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where $|\psi\rangle \in H = H_A \otimes H_B$ with $\dim H_A, \dim H_B < \infty$, $\{E_a^x\}$ PVMs on H_A for $\{F_b^y\}$ PVMs on H_B .
(finite dimensional separable quantum systems)
- **quantum spatial** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where we allow $\dim H_A = \dim H_B = \infty$.
(Infinite dimensional separable quantum systems)
- **quantum commuting** if $p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$,
where E_a^x, F_b^y PVMs on H and $[E_a^x, F_b^y] = 0$.
(Infinite dimensional and commuting observables)

local \subseteq quantum \subseteq quantum spatial \subseteq quantum commuting

Correlation under different assumptions

A correlation matrix $\{p(a, b|x, y)\}$ is called

- **local** if $p(a, b|x, y) = \sum_n \lambda_n p_n(a|x) q_n(b|y)$ ($\dim H = 1$, hidden variable)
- **quantum** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where $|\psi\rangle \in H = H_A \otimes H_B$ with $\dim H_A, \dim H_B < \infty$, $\{E_a^x\}$ PVMs on H_A for $\{F_b^y\}$ PVMs on H_B .
(finite dimensional separable quantum systems)
- **quantum spatial** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where we allow $\dim H_A = \dim H_B = \infty$.
(Infinite dimensional separable quantum systems)
- **quantum commuting** if $p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$,
where E_a^x, F_b^y PVMs on H and $[E_a^x, F_b^y] = 0$.
(Infinite dimensional and commuting observables)

local \subseteq quantum \subseteq quantum spatial \subseteq quantum commuting

Correlation under different assumptions

A correlation matrix $\{p(a, b|x, y)\}$ is called

- **local** if $p(a, b|x, y) = \sum_n \lambda_n p_n(a|x) q_n(b|y)$ ($\dim H = 1$, hidden variable)
- **quantum** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where $|\psi\rangle \in H = H_A \otimes H_B$ with $\dim H_A, \dim H_B < \infty$, $\{E_a^x\}$ PVMs on H_A for $\{F_b^y\}$ PVMs on H_B .
(finite dimensional separable quantum systems)
- **quantum spatial** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where we allow $\dim H_A = \dim H_B = \infty$.
(Infinite dimensional separable quantum systems)
- **quantum commuting** if $p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$,
where E_a^x, F_b^y PVMs on H and $[E_a^x, F_b^y] = 0$.
(Infinite dimensional and commuting observables)

local \subseteq quantum \subseteq quantum spatial \subseteq quantum commuting

Correlation under different assumptions

A correlation matrix $\{p(a, b|x, y)\}$ is called

- **local** if $p(a, b|x, y) = \sum_n \lambda_n p_n(a|x) q_n(b|y)$ ($\dim H = 1$, hidden variable)
- **quantum** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where $|\psi\rangle \in H = H_A \otimes H_B$ with $\dim H_A, \dim H_B < \infty$, $\{E_a^x\}$ PVMs on H_A for $\{F_b^y\}$ PVMs on H_B .
(finite dimensional **separable** quantum systems)
- **quantum spatial** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where we allow $\dim H_A = \dim H_B = \infty$.
(Infinite dimensional **separable** quantum systems)
- **quantum commuting** if $p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$,
where E_a^x, F_b^y PVMs on H and $[E_a^x, F_b^y] = 0$.
(Infinite dimensional and **commuting** observables)

local \subseteq quantum \subseteq quantum spatial \subseteq quantum commuting

Correlation under different assumptions

A correlation matrix $\{p(a, b|x, y)\}$ is called

- **local** if $p(a, b|x, y) = \sum_n \lambda_n p_n(a|x) q_n(b|y)$ ($\dim H = 1$, hidden variable)
- **quantum** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where $|\psi\rangle \in H = H_A \otimes H_B$ with $\dim H_A, \dim H_B < \infty$, $\{E_a^x\}$ PVMs on H_A for $\{F_b^y\}$ PVMs on H_B .
(finite dimensional **separable** quantum systems)
- **quantum spatial** if $p(a, b|x, y) = \langle \psi | E_a^x \otimes F_b^y | \psi \rangle$,
where we allow $\dim H_A = \dim H_B = \infty$.
(Infinite dimensional **separable** quantum systems)
- **quantum commuting** if $p(a, b|x, y) = \langle \psi | E_a^x F_b^y | \psi \rangle$,
where E_a^x, F_b^y PVMs on H and $[E_a^x, F_b^y] = 0$.
(Infinite dimensional and **commuting** observables)

local \subseteq quantum \subseteq quantum spatial \subseteq quantum commuting

Hierarchy of correlation sets

Let $C_{loc}, C_q, C_{qs}, C_{qc}$ be the set of respectively the *local*, *quantum*, *quantum spatial*, *quantum commuting* correlations set. For each (d, m) ,

$$C_{loc}(d, m) \subset C_q(d, m) \subset C_{qs}(d, m) \subset \overline{C_{qs}(d, m)} \subset C_{qc}(d, m) .$$

- $d = 1$: collapse to classic correlations, $C_{loc}(1, m) = C_{qc}(1, m)$
- $C_{loc}(2, 2) \neq C_q(2, 2) = C_{qc}(2, 2)$: Bell's inequality/CHSH game [Bell,'64] and experiments e.g. [Hanson et al,'15]

For $d, m \geq 2, (d, m) \neq (2, 2), C_{loc}(d, m) \neq C_q(d, m)$.

- Tsirelson's Problem '93: $\overline{C_{qs}(d, m)} \stackrel{?}{=} C_{qc}(d, m)$ for all $(d, m) > (2, 2)$. Equivalent to Connes Embedding Problem! [Junge et al,'11],[Fritz '11],[Ozawa, '13].
- [Slofstra, '17]: for $d \approx 100, C_{qs}(d, 8) \neq \overline{C_{qs}(d, 8)}$. Explicit examples. [Dykma-Paulsen-Prakash, '17]: $C_{qs}(5, 2) \neq \overline{C_{qs}(5, 2)}$. Existence proof.
- [Coladangelo-Stark, '18]: $C_q(5, 4) \neq C_{qs}(5, 4)$. Explicit examples.

Hierarchy of correlation sets

Let $C_{loc}, C_q, C_{qs}, C_{qc}$ be the set of respectively the *local*, *quantum*, *quantum spatial*, *quantum commuting* correlations set. For each (d, m) ,

$$C_{loc}(d, m) \subset C_q(d, m) \subset C_{qs}(d, m) \subset \overline{C_{qs}(d, m)} \subset C_{qc}(d, m) .$$

- $d = 1$: collapse to classic correlations, $C_{loc}(1, m) = C_{qc}(1, m)$
- $C_{loc}(2, 2) \neq C_q(2, 2) = C_{qc}(2, 2)$: Bell's inequality/CHSH game [Bell,'64] and experiments e.g. [Hanson et al,'15]

For $d, m \geq 2, (d, m) \neq (2, 2), C_{loc}(d, m) \neq C_q(d, m)$.

- Tsirelson's Problem '93: $\overline{C_{qs}(d, m)} \stackrel{?}{=} C_{qc}(d, m)$ for all $(d, m) > (2, 2)$. Equivalent to Connes Embedding Problem! [Junge et al,'11],[Fritz '11],[Ozawa, '13].
- [Slofstra, '17]: for $d \approx 100, C_{qs}(d, 8) \neq \overline{C_{qs}(d, 8)}$. Explicit examples. [Dykma-Paulsen-Prakash, '17]: $C_{qs}(5, 2) \neq \overline{C_{qs}(5, 2)}$. Existence proof.
- [Coladangelo-Stark, '18]: $C_q(5, 4) \neq C_{qs}(5, 4)$. Explicit examples.

Hierarchy of correlation sets

Let $C_{loc}, C_q, C_{qs}, C_{qc}$ be the set of respectively the *local*, *quantum*, *quantum spatial*, *quantum commuting* correlations set. For each (d, m) ,

$$C_{loc}(d, m) \subset C_q(d, m) \subset C_{qs}(d, m) \subset \overline{C_{qs}(d, m)} \subset C_{qc}(d, m) .$$

- $d = 1$: collapse to classic correlations, $C_{loc}(1, m) = C_{qc}(1, m)$
- $C_{loc}(2, 2) \neq C_q(2, 2) = C_{qc}(2, 2)$: Bell's inequality/CHSH game [Bell,'64] and experiments e.g. [Hanson et al,'15]

For $d, m \geq 2, (d, m) \neq (2, 2), C_{loc}(d, m) \neq C_q(d, m)$.

- Tsirelson's Problem '93: $\overline{C_{qs}(d, m)} \stackrel{?}{=} C_{qc}(d, m)$ for all $(d, m) > (2, 2)$. Equivalent to Connes Embedding Problem! [Junge et al,'11],[Fritz '11],[Ozawa, '13].
- [Slofstra, '17]: for $d \approx 100, C_{qs}(d, 8) \neq \overline{C_{qs}(d, 8)}$. Explicit examples. [Dykma-Paulsen-Prakash, '17]: $C_{qs}(5, 2) \neq \overline{C_{qs}(5, 2)}$. Existence proof.
- [Coladangelo-Stark, '18]: $C_q(5, 4) \neq C_{qs}(5, 4)$. Explicit examples.

Hierarchy of correlation sets

Let $C_{loc}, C_q, C_{qs}, C_{qc}$ be the set of respectively the *local*, *quantum*, *quantum spatial*, *quantum commuting* correlations set. For each (d, m) ,

$$C_{loc}(d, m) \subset C_q(d, m) \subset C_{qs}(d, m) \subset \overline{C_{qs}(d, m)} \subset C_{qc}(d, m) .$$

- $d = 1$: collapse to classic correlations, $C_{loc}(1, m) = C_{qc}(1, m)$
- $C_{loc}(2, 2) \neq C_q(2, 2) = C_{qc}(2, 2)$: Bell's inequality/CHSH game [Bell,'64] and experiments e.g. [Hanson et al,'15]

For $d, m \geq 2, (d, m) \neq (2, 2), C_{loc}(d, m) \neq C_q(d, m)$.

- Tsirelson's Problem '93: $\overline{C_{qs}(d, m)} \stackrel{?}{=} C_{qc}(d, m)$ for all $(d, m) > (2, 2)$. Equivalent to Connes Embedding Problem! [Junge et al,'11],[Fritz '11],[Ozawa, '13].
- [Slofstra, '17]: for $d \approx 100, C_{qs}(d, 8) \neq \overline{C_{qs}(d, 8)}$. Explicit examples. [Dykma-Paulsen-Prakash, '17]: $C_{qs}(5, 2) \neq \overline{C_{qs}(5, 2)}$. Existence proof.
- [Coladangelo-Stark, '18]: $C_q(5, 4) \neq C_{qs}(5, 4)$. Explicit examples.

Hierarchy of correlation sets

Let $C_{loc}, C_q, C_{qs}, C_{qc}$ be the set of respectively the *local*, *quantum*, *quantum spatial*, *quantum commuting* correlations set. For each (d, m) ,

$$C_{loc}(d, m) \subset C_q(d, m) \subset C_{qs}(d, m) \subset \overline{C_{qs}(d, m)} \subset C_{qc}(d, m) .$$

- $d = 1$: collapse to classic correlations, $C_{loc}(1, m) = C_{qc}(1, m)$
- $C_{loc}(2, 2) \neq C_q(2, 2) = C_{qc}(2, 2)$: Bell's inequality/CHSH game [Bell,'64] and experiments e.g. [Hanson et al,'15]

For $d, m \geq 2, (d, m) \neq (2, 2), C_{loc}(d, m) \neq C_q(d, m)$.

- Tsirelson's Problem '93: $\overline{C_{qs}(d, m)} \stackrel{?}{=} C_{qc}(d, m)$ for all $(d, m) > (2, 2)$. Equivalent to Connes Embedding Problem! [Junge et al,'11],[Fritz '11],[Ozawa, '13].
- [Slofstra, '17]: for $d \approx 100, C_{qs}(d, 8) \neq \overline{C_{qs}(d, 8)}$. Explicit examples. [Dykma-Paulsen-Prakash, '17]: $C_{qs}(5, 2) \neq \overline{C_{qs}(5, 2)}$. Existence proof.
- [Coladangelo-Stark, '18]: $C_q(5, 4) \neq C_{qs}(5, 4)$. Explicit examples.

Hierarchy of correlation sets

Let $C_{loc}, C_q, C_{qs}, C_{qc}$ be the set of respectively the *local*, *quantum*, *quantum spatial*, *quantum commuting* correlations set. For each (d, m) ,

$$C_{loc}(d, m) \subset C_q(d, m) \subset C_{qs}(d, m) \subset \overline{C_{qs}(d, m)} \subset C_{qc}(d, m).$$

- $d = 1$: collapse to classic correlations, $C_{loc}(1, m) = C_{qc}(1, m)$
- $C_{loc}(2, 2) \neq C_q(2, 2) = C_{qc}(2, 2)$: Bell's inequality/CHSH game [Bell,'64] and experiments e.g. [Hanson et al,'15]

For $d, m \geq 2, (d, m) \neq (2, 2), C_{loc}(d, m) \neq C_q(d, m)$.

- Tsirelson's Problem '93: $\overline{C_{qs}(d, m)} \stackrel{?}{=} C_{qc}(d, m)$ for all $(d, m) > (2, 2)$. Equivalent to Connes Embedding Problem! [Junge et al,'11],[Fritz '11],[Ozawa, '13].
- [Slofstra, '17]: for $d \approx 100, C_{qs}(d, 8) \neq \overline{C_{qs}(d, 8)}$. Explicit examples. [Dykma-Paulsen-Prakash, '17]: $C_{qs}(5, 2) \neq \overline{C_{qs}(5, 2)}$. Existence proof.
- [Coladangelo-Stark, '18]: $C_q(5, 4) \neq C_{qs}(5, 4)$. Explicit examples.

C_q vs. C_{qs} vs. $\overline{C_{qs}}$ at smallest possible sizes?

$C_q(3,2) \neq C_{qs}(3,2) \neq \overline{C_{qs}(3,2)}$?, $C_q(2,3) \neq C_{qs}(2,3) \neq \overline{C_{qs}(2,3)}$?

Matrix valued quantum correlations

We consider the following matrix-valued quantum correlations:

$$C_{qs}^n(d, m) = \left\{ \left[V^*(E_a^x \otimes F_b^y)V \right]_{a,b}^{x,y} \mid \begin{array}{l} H_A, H_B \text{ Hilbert spaces,} \\ V : \mathbb{C}^n \rightarrow H_A \otimes H_B \text{ an isometry,} \\ (E_a^x)_{x=1}^m \text{ PVMs on } H_A \text{ for each } a \\ (F_b^y)_{y=1}^m \text{ PVMs on } H_B \text{ for each } b \end{array} \right\}$$

$$C_{qc}^n(d, m) = [E_a^x, F_b^y] = 0 \text{ on one Hilbert space } H$$

$$C_q^n(d, m) = H_A \text{ and } H_B \text{ finite dimensional}$$

Quantum correlations set $C_q(d, m) = C_q^1(d, m)$

$C_q^n(d, m)$: a vector $\psi \rightsquigarrow n$ -dimensional subspace. Corresponds to games requires quantum answers to classical questions.

- [Junge et al, '11],[Fritz '11]: For any fixed $(d, m) > (2, 2)$,

$$\overline{C_{qs}^n(d, m)} = C_{qc}^n(d, m) \quad \forall n \geq 1 \Leftrightarrow \text{CEC}$$

- [Ozawa, '13]:

$$\text{CEC} \Leftrightarrow \overline{C_{qs}(d, m)} = C_{qc}(d, m) \quad \forall (d, m)$$

Matrix valued quantum correlations

We consider the following matrix-valued quantum correlations:

$$C_{qs}^n(d, m) = \left\{ \left[V^*(E_a^x \otimes F_b^y)V \right]_{a,b}^{x,y} \mid \begin{array}{l} H_A, H_B \text{ Hilbert spaces,} \\ V : \mathbb{C}^n \rightarrow H_A \otimes H_B \text{ an isometry,} \\ (E_a^x)_{x=1}^m \text{ PVMs on } H_A \text{ for each } a \\ (F_b^y)_{y=1}^m \text{ PVMs on } H_B \text{ for each } b \end{array} \right\}$$

$$C_{qc}^n(d, m) = [E_a^x, F_b^y] = 0 \text{ on one Hilbert space } H$$

$$C_q^n(d, m) = H_A \text{ and } H_B \text{ finite dimensional}$$

Quantum correlations set $C_q(d, m) = C_q^1(d, m)$

$C_q^n(d, m)$: a vector $\psi \rightsquigarrow n$ -dimensional subspace. Corresponds to games requires quantum answers to classical questions.

- [Junge et al, '11],[Fritz '11]: For any fixed $(d, m) > (2, 2)$,

$$\overline{C_{qs}^n(d, m)} = C_{qc}^n(d, m) \forall n \geq 1 \Leftrightarrow \text{CEC}$$

- [Ozawa, '13]:

$$\text{CEC} \Leftrightarrow \overline{C_{qs}(d, m)} = C_{qc}(d, m) \forall (d, m)$$

Matrix valued quantum correlations

We consider the following matrix-valued quantum correlations:

$$C_{qs}^n(d, m) = \left\{ \begin{array}{l} \left[V^* (E_a^x \otimes F_b^y) V \right]_{a,b}^{x,y} \mid \begin{array}{l} H_A, H_B \text{ Hilbert spaces,} \\ V : \mathbb{C}^n \rightarrow H_A \otimes H_B \text{ an isometry,} \\ (E_a^x)_{x=1}^m \text{ PVMs on } H_A \text{ for each } a \\ (F_b^y)_{y=1}^m \text{ PVMs on } H_B \text{ for each } b \end{array} \end{array} \right\}$$

$$C_{qc}^n(d, m) = [E_a^x, F_b^y] = 0 \text{ on one Hilbert space } H$$

$$C_q^n(d, m) = H_A \text{ and } H_B \text{ finite dimensional}$$

Quantum correlations set $C_q(d, m) = C_q^1(d, m)$

$C_q^n(d, m)$: a vector $\psi \rightsquigarrow n$ -dimensional subspace. Corresponds to games requires quantum answers to classical questions.

- [Junge et al, '11], [Fritz '11]: For any fixed $(d, m) > (2, 2)$,

$$\overline{C_{qs}^n(d, m)} = C_{qc}^n(d, m) \quad \forall n \geq 1 \Leftrightarrow \text{CEC}$$

- [Ozawa, '13]:

$$\text{CEC} \Leftrightarrow \overline{C_{qs}(d, m)} = C_{qc}(d, m) \quad \forall (d, m)$$

Theorem (G.-Harris-Junge, 2017)

$$C_{qs}^3(4, 2) \neq \overline{C_{qs}^3(4, 2)}, C_{qs}^5(3, 2) \neq \overline{C_{qs}^5(3, 2)}, C_{qs}^{13}(2, 3) \neq \overline{C_{qs}^{13}(2, 3)}$$

An infinite dimensional correlation

Lemma

Let H be a Hilbert space. There exist unitaries u_0, u_1, v_0, v_1 on H and orthonormal vectors $|h_0\rangle, |h_1\rangle \in H \otimes H$ satisfying

$$u_0 \otimes v_0 |h_0\rangle = u_1 \otimes v_1 |h_0\rangle = (|h_0\rangle + |h_1\rangle)/\sqrt{2},$$

$$u_0 \otimes v_1 |h_0\rangle = u_1 \otimes v_0 |h_0\rangle = (|h_0\rangle - |h_1\rangle)/\sqrt{2},$$

if and only if H is infinite dimensional.

Proof. Denote $u := u_1^* u_0, v := v_1^* v_0$.

$$u \otimes 1 |h_0\rangle = 1 \otimes v |h_0\rangle = |h_0\rangle, u \otimes 1 |h_1\rangle = 1 \otimes v |h_1\rangle = -|h_0\rangle.$$

$$|h_0\rangle \in H_+ \otimes K_+, |h_1\rangle \in H_- \otimes K_-$$

where H_+ (resp. K_+) is the eigenspace of u (resp. v) of eigenvalue 1, H_- (resp. K_-) is the eigenspace of u (resp. v) of eigenvalue -1 .

$H_+ \perp H_-, K_+ \perp K_-$, hence for Schmidt Rank "Sr",

$$Sr(|h_0\rangle + |h_1\rangle) = Sr(|h_0\rangle) + Sr(|h_1\rangle), Sr(|h_0\rangle - |h_1\rangle) = Sr(|h_0\rangle).$$

$$\Rightarrow Sr(|h_0\rangle) = \infty.$$

An infinite dimensional correlation

Lemma

Let H be a Hilbert space. There exist unitaries u_0, u_1, v_0, v_1 on H and orthonormal vectors $|h_0\rangle, |h_1\rangle \in H \otimes H$ satisfying

$$u_0 \otimes v_0 |h_0\rangle = u_1 \otimes v_1 |h_0\rangle = (|h_0\rangle + |h_1\rangle)/\sqrt{2},$$

$$u_0 \otimes v_1 |h_0\rangle = u_1 \otimes v_0 |h_0\rangle = (|h_0\rangle - |h_1\rangle)/\sqrt{2},$$

if and only if H is infinite dimensional.

Proof. Denote $u := u_1^* u_0, v := v_1^* v_0$.

$$u \otimes 1(|h_0\rangle + |h_1\rangle) = 1 \otimes v(|h_0\rangle + |h_1\rangle) = (|h_0\rangle - |h_1\rangle),$$

$$u \otimes v(|h_0\rangle + |h_1\rangle) = (|h_0\rangle + |h_1\rangle)$$

$$u \otimes 1|h_0\rangle = 1 \otimes v|h_0\rangle = |h_0\rangle, u \otimes 1|h_1\rangle = 1 \otimes v|h_1\rangle = -|h_0\rangle.$$

$$|h_0\rangle \in H_+ \otimes K_+, |h_1\rangle \in H_- \otimes K_-$$

where H_+ (resp. K_+) is the eigenspace of u (resp. v) of eigenvalue 1, H_- (resp. K_-) is the eigenspace of u (resp. v) of eigenvalue -1.

An infinite dimensional correlation

Lemma

Let H be a Hilbert space. There exist unitaries u_0, u_1, v_0, v_1 on H and orthonormal vectors $|h_0\rangle, |h_1\rangle \in H \otimes H$ satisfying

$$u_0 \otimes v_0 |h_0\rangle = u_1 \otimes v_1 |h_0\rangle = (|h_0\rangle + |h_1\rangle)/\sqrt{2},$$

$$u_0 \otimes v_1 |h_0\rangle = u_1 \otimes v_0 |h_0\rangle = (|h_0\rangle - |h_1\rangle)/\sqrt{2},$$

if and only if H is infinite dimensional.

Proof. Denote $u := u_1^* u_0, v := v_1^* v_0$.

$$u \otimes 1 |h_0\rangle = 1 \otimes v |h_0\rangle = |h_0\rangle, u \otimes 1 |h_1\rangle = 1 \otimes v |h_1\rangle = -|h_0\rangle.$$

$$|h_0\rangle \in H_+ \otimes K_+, |h_1\rangle \in H_- \otimes K_-$$

where H_+ (resp. K_+) is the eigenspace of u (resp. v) of eigenvalue 1, H_- (resp. K_-) is the eigenspace of u (resp. v) of eigenvalue -1 .

$H_+ \perp H_-, K_+ \perp K_-$, hence for Schmidt Rank "Sr",

$$Sr(|h_0\rangle + |h_1\rangle) = Sr(|h_0\rangle) + Sr(|h_1\rangle), Sr(|h_0\rangle - |h_1\rangle) = Sr(|h_0\rangle).$$

$$\Rightarrow Sr(|h_0\rangle) = \infty.$$

An infinite dimensional correlation

Lemma

Let H be a Hilbert space. There exist unitaries u_0, u_1, v_0, v_1 on H and orthonormal vectors $|h_0\rangle, |h_1\rangle \in H \otimes H$ satisfying

$$u_0 \otimes v_0 |h_0\rangle = u_1 \otimes v_1 |h_0\rangle = (|h_0\rangle + |h_1\rangle)/\sqrt{2},$$

$$u_0 \otimes v_1 |h_0\rangle = u_1 \otimes v_0 |h_0\rangle = (|h_0\rangle - |h_1\rangle)/\sqrt{2},$$

if and only if H is infinite dimensional.

Proof. Denote $u := u_1^* u_0, v := v_1^* v_0$.

$$u \otimes 1 |h_0\rangle = 1 \otimes v |h_0\rangle = |h_0\rangle, u \otimes 1 |h_1\rangle = 1 \otimes v |h_1\rangle = -|h_1\rangle.$$

$$|h_0\rangle \in H_+ \otimes K_+, |h_1\rangle \in H_- \otimes K_-$$

where H_+ (resp. K_+) is the eigenspace of u (resp. v) of eigenvalue 1, H_- (resp. K_-) is the eigenspace of u (resp. v) of eigenvalue -1 .

$H_+ \perp H_-, K_+ \perp K_-$, hence for Schmidt Rank “ Sr ”,

$$Sr(|h_0\rangle + |h_1\rangle) = Sr(|h_0\rangle) + Sr(|h_1\rangle), Sr(|h_0\rangle - |h_1\rangle) = Sr(|h_0\rangle).$$

$$\Rightarrow Sr(|h_0\rangle) = \infty.$$

A non-spatial correlation

Lemma

There are no unitary $u_0, u_1, u_2, v_0, v_1, v_2$ on H and two orthonormal vector $|h_0\rangle, |h_1\rangle \in H \otimes H$ satisfying the equations

$$u_0 \otimes v_0 |h_0\rangle = u_1 \otimes v_1 |h_0\rangle = (|h_0\rangle + |h_1\rangle)\sqrt{2}, \quad (1)$$

$$u_0 \otimes v_1 |h_0\rangle = u_1 \otimes v_0 |h_0\rangle = (|h_0\rangle - |h_1\rangle)\sqrt{2}, \quad (2)$$

$$u_2 \otimes v_2 |h_0\rangle = |h_1\rangle. \quad (3)$$

Proof. (1), (2) implies

$$\begin{aligned} & \text{Schmidt coefficients}\left(\frac{|h_0\rangle + |h_1\rangle}{\sqrt{2}}\right) \\ &= \frac{1}{\sqrt{2}} \text{Schmidt coefficients}(|h_0\rangle) \cup \frac{1}{\sqrt{2}} \text{Schmidt coefficients}(|h_1\rangle). \end{aligned}$$

(3) implies $|h_0\rangle, |h_1\rangle$ and $\frac{1}{\sqrt{2}}(|h_0\rangle + |h_1\rangle)$ have the same Schmidt coefficients, which is a contradiction.

Approximation by Embezzlement of Entanglement

Theorem (van Dam-Hayden,'02)

For any $\epsilon > 0$, there exists a finite dimensional H , a unit vector $|R\rangle \in H \otimes H$ and unitaries U, V on $\mathbb{C}^2 \otimes H$ such that

$$\|U \otimes V|0\rangle|R\rangle|0\rangle - \frac{1}{\sqrt{2}}(|0\rangle|R\rangle|0\rangle + |1\rangle|R\rangle|1\rangle)\| \leq \epsilon.$$

Let $X = \begin{bmatrix} 1 & \\ & -1 \end{bmatrix}$, $Z = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$. Choose

$$u_0 = U, v_0 = V, u_1 = (X \otimes 1)U, v_1 = (X \otimes 1)V,$$

$$u_2 = v_2 = Z \otimes 1,$$

$$|h_0\rangle = |0\rangle|R\rangle|0\rangle, |h_1\rangle = |1\rangle|R\rangle|1\rangle.$$

When $\epsilon \rightarrow 0$, $u_0 \otimes v_0|h_0\rangle \sim \frac{|h_0\rangle + |h_1\rangle}{\sqrt{2}} \sim u_1 \otimes v_1|h_0\rangle$

$u_0 \otimes v_1|h_0\rangle \sim \frac{|h_0\rangle - |h_1\rangle}{\sqrt{2}} \sim u_1 \otimes v_0|h_0\rangle$

Approximation by Embezzlement of Entanglement

Theorem (van Dam-Hayden,'02)

For any $\epsilon > 0$, there exists a finite dimensional H , a unit vector $|R\rangle \in H \otimes H$ and unitaries U, V on $\mathbb{C}^2 \otimes H$ such that

$$\|U \otimes V|0\rangle|R\rangle|0\rangle - \frac{1}{\sqrt{2}}(|0\rangle|R\rangle|0\rangle + |1\rangle|R\rangle|1\rangle)\| \leq \epsilon.$$

Let $X = \begin{bmatrix} 1 & \\ & -1 \end{bmatrix}$, $Z = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$. Choose

$$u_0 = U, v_0 = V, u_1 = (X \otimes 1)U, v_1 = (X \otimes 1)V,$$

$$u_2 = v_2 = Z \otimes 1,$$

$$|h_0\rangle = |0\rangle|R\rangle|0\rangle, |h_1\rangle = |1\rangle|R\rangle|1\rangle.$$

When $\epsilon \rightarrow 0$, $u_0 \otimes v_0|h_0\rangle \sim \frac{|h_0\rangle + |h_1\rangle}{\sqrt{2}} \sim u_1 \otimes v_1|h_0\rangle$

$u_0 \otimes v_1|h_0\rangle \sim \frac{|h_0\rangle - |h_1\rangle}{\sqrt{2}} \sim u_1 \otimes v_0|h_0\rangle$

Matrix-valued correlation by group embeddings

- n -valued PVM $\sum_{j=0}^{n-1} E_j = 1 \Leftrightarrow$ order n unitary $u = \sum_{j=0}^{n-1} e^{\frac{2\pi j i}{n}} E_j$.
- In previous lemma, we have two order 2 unitary X, Z and an abstract unitary U in the limit $\epsilon \rightarrow 0$ has infinity order, hence corresponds to a vector state of the unitary representation of

$$\mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z} \times \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}$$

- Four 2-valued quantum correlations \rightsquigarrow states of unitary representation of

$$\mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 \times \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2$$

- Using group embedding, $\mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z} \hookrightarrow \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2$

$$X \mapsto \sigma_0, Z \mapsto \sigma_1, U \mapsto \sigma_2 \sigma_3.$$

- We have a non-spatial state witnessed on the set

$$\{\sigma_0, \sigma_1, \sigma_2 \sigma_3, \sigma_1 \sigma_2 \sigma_3\} \otimes \{\sigma_0, \sigma_1, \sigma_2 \sigma_3, \sigma_1 \sigma_2 \sigma_3\}$$

Matrix-valued correlation by group embeddings

- n -valued PVM $\sum_{j=0}^{n-1} E_j = 1 \Leftrightarrow$ order n unitary $u = \sum_{j=0}^{n-1} e^{\frac{2\pi j i}{n}} E_j$.
- In previous lemma, we have two order 2 unitary X, Z and an abstract unitary U in the limit $\epsilon \rightarrow 0$ has infinity order, hence corresponds to a vector state of the unitary representation of

$$\mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z} \times \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}$$

- Four 2-valued quantum correlations \rightsquigarrow states of unitary representation of

$$\mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 \times \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2$$

- Using group embedding, $\mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z} \hookrightarrow \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2$

$$X \mapsto \sigma_0, Z \mapsto \sigma_1, U \mapsto \sigma_2 \sigma_3.$$

- We have a non-spatial state witnessed on the set

$$\{\sigma_0, \sigma_1, \sigma_2 \sigma_3, \sigma_1 \sigma_2 \sigma_3\} \otimes \{\sigma_0, \sigma_1, \sigma_2 \sigma_3, \sigma_1 \sigma_2 \sigma_3\}$$

Matrix-valued correlation by group embeddings

- n -valued PVM $\sum_{j=0}^{n-1} E_j = 1 \Leftrightarrow$ order n unitary $u = \sum_{j=0}^{n-1} e^{\frac{2\pi j i}{n}} E_j$.
- In previous lemma, we have two order 2 unitary X, Z and an abstract unitary U in the limit $\epsilon \rightarrow 0$ has infinity order, hence corresponds to a vector state of the unitary representation of

$$\mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z} \times \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}$$

- Four 2-valued quantum correlations \rightsquigarrow states of unitary representation of

$$\mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 \times \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2$$

- Using group embedding, $\mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z} \hookrightarrow \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2 * \mathbb{Z}_2$

$$X \mapsto \sigma_0, Z \mapsto \sigma_1, U \mapsto \sigma_2 \sigma_3.$$

- We have a non-spatial state witnessed on the set

$$\{\sigma_0, \sigma_1, \sigma_2 \sigma_3, \sigma_1 \sigma_2 \sigma_3\} \otimes \{\sigma_0, \sigma_1, \sigma_2 \sigma_3, \sigma_1 \sigma_2 \sigma_3\}$$

A concrete example in $\overline{C_{qs}^3(4, 2)}$

Let $\sigma_0, \sigma_1, \sigma_2, \sigma_3$ be self adjoint unitaries. $\sigma_x = E_x^+ - E_x^-$

$$\Psi(\sigma_0 \otimes \sigma_0) = \begin{bmatrix} 0 & 1 & * \\ 1 & 0 & * \\ * & * & * \end{bmatrix}, \quad \Psi(\sigma_2 \otimes \sigma_2) = \frac{1}{\sqrt{2}} \begin{bmatrix} * & * & 1 \\ * & * & 1 \\ 1 & 1 & * \end{bmatrix},$$

$$\Psi(\sigma_3 \otimes \sigma_3) = \begin{bmatrix} * & * & 1 \\ * & * & * \\ 1 & * & * \end{bmatrix}, \quad \Psi(\sigma_1 \otimes \sigma_1) = \begin{bmatrix} 1 & 0 & * \\ 0 & 1 & * \\ * & * & * \end{bmatrix},$$

$$\Psi(\sigma_1 \otimes 1) = \begin{bmatrix} 1 & 0 & * \\ 0 & -1 & * \\ * & * & * \end{bmatrix}, \quad \Psi(1 \otimes \sigma_1) = \begin{bmatrix} 1 & 0 & * \\ 0 & -1 & * \\ * & * & * \end{bmatrix},$$

Separation between C_q^n vs C_{qs}^n .

Theorem (G.-Harris-Junge, 2017)

$$C_{qs}^3(4, 2) \neq \overline{C_{qs}^3(4, 2)}, C_{qs}^5(3, 2) \neq \overline{C_{qs}^5(3, 2)}, C_{qs}^{13}(2, 3) \neq \overline{C_{qs}^{13}(2, 3)}$$

Theorem (Harris, 2018)

$$C_q^2(3, 2) \neq C_{qs}^2(3, 2), C_q^4(2, 3) \neq C_{qs}^4(3, 2).$$

Observation:

C_{qs} vs. $\overline{C_{qs}}$: infinite “amount” of entanglement (entropy)

C_q vs. C_{qs} : infinite “rank” of entanglement (Schmidt rank)

Separation between C_q^n vs C_{qs}^n .

Theorem (G.-Harris-Junge, 2017)

$$C_{qs}^3(4, 2) \neq \overline{C_{qs}^3(4, 2)}, C_{qs}^5(3, 2) \neq \overline{C_{qs}^5(3, 2)}, C_{qs}^{13}(2, 3) \neq \overline{C_{qs}^{13}(2, 3)}$$

Theorem (Harris, 2018)

$$C_q^2(3, 2) \neq C_{qs}^2(3, 2), C_q^4(2, 3) \neq C_{qs}^4(3, 2).$$

Observation:

C_{qs} vs. $\overline{C_{qs}}$: infinite “amount” of entanglement (entropy)

C_q vs. C_{qs} : infinite “rank” of entanglement (Schmidt rank)

Perfect embezzlement of entanglement

Perfect embezzlement of entanglement: a unit vector $|R\rangle \in H \otimes H$ and unitaries U, V on $\mathbb{C}^2 \otimes H$ such that

$$U \otimes V|0\rangle|R\rangle|0\rangle = \frac{1}{\sqrt{2}}(|0\rangle|R\rangle|0\rangle + |1\rangle|R\rangle|1\rangle).$$

Let $U = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}$ and $V = \begin{bmatrix} v_{00} & v_{01} \\ v_{10} & v_{11} \end{bmatrix}$. This means

$$\langle R|u_{00} \otimes v_{00}|R\rangle = \langle R|u_{10} \otimes v_{10}|R\rangle = \frac{1}{\sqrt{2}},$$

$$\langle R|u_{10} \otimes v_{00}|R\rangle = \langle R|u_{00} \otimes v_{10}|R\rangle = 0.$$

Teleportation and dense coding for unitary

2×2 unitary $U = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix} \xrightarrow{\mathcal{T}} 4$ unitaries

$$U = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}, (1 \otimes Z)U(1 \otimes Z) = \begin{bmatrix} u_{00} & -u_{01} \\ -u_{10} & u_{11} \end{bmatrix},$$

$$(1 \otimes X)U(1 \otimes X) = \begin{bmatrix} u_{11} & u_{10} \\ u_{01} & u_{00} \end{bmatrix}, (1 \otimes Y)U(1 \otimes Y) = \begin{bmatrix} u_{11} & -u_{10} \\ -u_{01} & u_{00} \end{bmatrix}$$

Teleportation and dense coding for unitary

2×2 unitary $U = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix} \xrightarrow{\mathcal{T}} 4$ unitaries

$$U = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}, (1 \otimes Z)U(1 \otimes Z) = \begin{bmatrix} u_{00} & -u_{01} \\ -u_{10} & u_{11} \end{bmatrix},$$

$$(1 \otimes X)U(1 \otimes X) = \begin{bmatrix} u_{11} & u_{10} \\ u_{01} & u_{00} \end{bmatrix}, (1 \otimes Y)U(1 \otimes Y) = \begin{bmatrix} u_{11} & -u_{10} \\ -u_{01} & u_{00} \end{bmatrix}$$

4 unitaries $g_1, g_2, g_3, g_4 \xrightarrow{\mathcal{S}} 2 \times 2$ unitary entries :

$$u_{00} = \frac{1}{2} \begin{bmatrix} g_1 + g_2 & \\ & g_3 + g_4 \end{bmatrix}, u_{01} = \frac{1}{2} \begin{bmatrix} & g_1 - g_2 \\ g_3 - g_4 & \end{bmatrix},$$

$$u_{10} = \frac{1}{2} \begin{bmatrix} & g_3 - g_4 \\ g_1 - g_2 & \end{bmatrix}, u_{11} = \frac{1}{2} \begin{bmatrix} g_3 + g_4 & \\ & g_1 + g_2 \end{bmatrix},$$

Teleportation and dense coding for unitary

2×2 unitary $U = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix} \xrightarrow{\mathcal{T}} 4$ unitaries

$$U = \begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}, (1 \otimes Z)U(1 \otimes Z) = \begin{bmatrix} u_{00} & -u_{01} \\ -u_{10} & u_{11} \end{bmatrix},$$

$$(1 \otimes X)U(1 \otimes X) = \begin{bmatrix} u_{11} & u_{10} \\ u_{01} & u_{00} \end{bmatrix}, (1 \otimes Y)U(1 \otimes Y) = \begin{bmatrix} u_{11} & -u_{10} \\ -u_{01} & u_{00} \end{bmatrix}$$

4 unitaries $g_1, g_2, g_3, g_4 \xrightarrow{\mathcal{S}} 2 \times 2$ unitary entries:

$$U = \frac{1}{2} \begin{bmatrix} g_1 + g_2 & & & g_1 - g_2 \\ & g_3 + g_4 & g_3 - g_4 & \\ & g_3 - g_4 & g_3 + g_4 & \\ g_1 - g_2 & & & g_1 + g_2 \end{bmatrix}$$

C^* -Teleportation and Dense Coding

$C^*(\mathbb{F}_4)$: the “universal” C^* -algebra of four unitaries $g_{00}, g_{10}, g_{01}, g_{11}$.

\mathcal{B}_2 : the “universal” C^* -algebra of $u_{00}, u_{01}, u_{10}, u_{11}$ such that $\begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}$ is a unitary.

Theorem (G.-Harris-Junge, 2017)

The teleportation and dense coding scheme gives the following C^* -algebras embedding:

$$B_2 \xrightarrow{\mathcal{T}} M_2(C^*(\mathbb{F}_4)) \longrightarrow B_2, \quad C^*(\mathbb{F}_4) \xrightarrow{\mathcal{S}} M_2(B_2) \longrightarrow C^*(\mathbb{F}_4)$$

$$\begin{array}{ccc} B_2 \otimes B_2 & \xrightarrow{\mathcal{T} \otimes \mathcal{T}} M_2(C^*(\mathbb{F}_4)) \otimes M_2(C^*(\mathbb{F}_4)) & \xrightarrow{\mathcal{S}^{(2)} \otimes \mathcal{S}^{(2)}} M_4(\mathcal{B}_2) \otimes M_4(\mathcal{B}_2) \\ \downarrow \psi & \nearrow \phi & \downarrow \simeq \\ \mathbb{C} & \xleftarrow{\rho \otimes \psi} & M_4 \otimes M_4 \otimes (\mathcal{B}_2 \otimes \mathcal{B}_2) \end{array}$$

ψ is the state of perfect embezzlement. ρ is an entangled bit.

ϕ is our M_2 -valued non-spatial correlation for unitaries.

C^* -Teleportation and Dense Coding

$C^*(\mathbb{F}_4)$: the “universal” C^* -algebra of four unitaries $g_{00}, g_{10}, g_{01}, g_{11}$.

\mathcal{B}_2 : the “universal” C^* -algebra of $u_{00}, u_{01}, u_{10}, u_{11}$ such that $\begin{bmatrix} u_{00} & u_{01} \\ u_{10} & u_{11} \end{bmatrix}$ is a unitary.

Theorem (G.-Harris-Junge, 2017)

The teleportation and dense coding scheme gives the following C^* -algebras embedding:

$$B_2 \xrightarrow{\mathcal{T}} M_2(C^*(\mathbb{F}_4)) \longrightarrow B_2, \quad C^*(\mathbb{F}_4) \xrightarrow{\mathcal{S}} M_2(B_2) \longrightarrow C^*(\mathbb{F}_4)$$

$$\begin{array}{ccc} B_2 \otimes B_2 & \xrightarrow{\mathcal{T} \otimes \mathcal{T}} M_2(C^*(\mathbb{F}_4)) \otimes M_2(C^*(\mathbb{F}_4)) & \xrightarrow{\mathcal{S}^{(2)} \otimes \mathcal{S}^{(2)}} M_4(\mathcal{B}_2) \otimes M_4(\mathcal{B}_2) \\ \psi \downarrow & \nearrow \phi & \downarrow \simeq \\ \mathbb{C} & \xleftarrow{\rho \otimes \psi} & M_4 \otimes M_4 \otimes (\mathcal{B}_2 \otimes \mathcal{B}_2) \end{array}$$

ψ is the state of perfect embezzlement. ρ is an entangled bit.

ϕ is our M_2 -valued non-spatial correlation for unitaries.

All Teleportation and Super Dense Coding Schemes

Let ϕ be a maximally entangled state.

Theorem (Werner, 2001)

There are one to one correspondence between the following objects:

- 1 *Super dense coding Schemes* $|j\rangle\langle j| \rightarrow |\phi_j\rangle = U_j \otimes 1|\phi\rangle \rightarrow |j\rangle\langle j|$
- 2 *Orthonormal bases of Maximally entangled states.* $|\phi_j\rangle = U_j \otimes 1|\phi\rangle$
- 3 *Unitary orthonormal bases in Hilbert-Schmidt norm i.e.* $\text{tr}(U_k^* U_j) = \delta_{jk}$
- 4 *Teleportation schemes* $\rho \rightarrow \frac{1}{d^2} \sum_{j=1} U_j \rho U_j^* \otimes |j\rangle\langle j| \rightarrow \rho$

Theorem (G.-Werner, 2017)

Teleportation & dense coding schemes are one to one corresponds

i) *completely positive 1-complemented embeddings to*

$$S_1^d \xrightarrow{\mathcal{T}} M_d(l_1^{d^2}) \longrightarrow S_1^d, \quad l_1^{d^2} \xrightarrow{\mathcal{S}} M_d(S_1^d) \longrightarrow l_1^{d^2}.$$

ii) *injective *-homomorphism which sends $\{u_{jk}\}$ to $M_d(\text{span}\{g_{jk}\})$*

$$B_d \xrightarrow{\mathcal{T}} M_d(C^*(\mathbb{F}_{d^2})) \longrightarrow B_d, \quad C^*(\mathbb{F}_{d^2}) \xrightarrow{\mathcal{S}} M_d(B_d) \longrightarrow C^*(\mathbb{F}_{d^2})$$

Summary

- 1 We have separation of Matrix valued quantum correlation for C_q^n vs. C_{qs}^n vs. $\overline{C_{qs}^n}$ at smallest measurement size $(d, m) = (3, 2), (2, 3)$.
- 2 Our correlations are explicit.
- 3 An interesting key lemma relates to embezzlement of entanglement. A quantum game scenario?

Remains Open

$$C_q(3, 2) \neq C_{qs}(3, 2) \neq \overline{C_{qs}(3, 2)}?, \quad C_q(2, 3) \neq C_{qs}(2, 3) \neq \overline{C_{qs}(2, 3)}?$$

Thank you for your attention!