Ultrafast 1 MHz vacuum-ultraviolet source via highly cascaded harmonic generation in negative-curvature hollow-core fibers

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Vacuum-ultraviolet (VUV) light is critical for the study of molecules and materials, but the generation of femtosecond pulses in the VUV region at high repetition rates has proven difficult. Here we demonstrate the efficient generation of VUV light at megahertz repetition rates using highly cascaded four-wave mixing processes in a negative-curvature hollow-core fiber. Both even- and odd-order harmonics are generated up to the 15th harmonic (69 nm, 18.0 eV), with high energy resolution of ~40 meV. In contrast to direct high harmonic generation, this highly cascaded harmonic generation process requires lower peak intensity and therefore can operate at higher repetition rates, driven by a robust 10 W fiber-laser system in a compact setup. Additionally, we present numerical simulations that explore the fundamental capabilities and spatiotemporal dynamics of highly cascaded harmonic generation. This VUV source can enhance the capabilities of spectroscopies of molecular and quantum materials, such as photoionization mass spectrometry and time-, angle-, and spin-resolved photoemission.

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1. INTRODUCTION

The vacuum-ultraviolet (VUV) spectral region, covering approximately 6–15 eV, has a unique ability to probe physical and chemical processes. For example, the bond energies of all molecules and the ionization energies of most materials lie in this energy range. Consequently, VUV light sources are used as an ionization source in angle-resolved photoelectron spectroscopy (ARPES) [1–5] and photoionization mass spectrometry (PIMS) [6–12], or to initiate chemical reactions relevant to atmospheric science in a controlled environment [13–15]. However, this science is limited by the lack of bright VUV light sources, especially coherent (laser) sources. While low-flux sources, such as deuterium lamps, can satisfy some applications [16], many experiments require higher flux to overcome shot noise in a reasonable amount of time. Single-wavelength sources, such as the ninth harmonic of a Nd:YAG laser (118.2 nm, 10.49 eV) [17], are not tunable, limiting the scope of experiments for which they can be used. Tunable deep-ultraviolet light has been produced from phase-matched four-wave mixing schemes [18,19], but extending this strategy into the VUV has typically provided very limited or no tunability [20–22]. Synchrotrons and free-electron lasers are currently the only sources of fully tunable, high-flux VUV light [8]; however, these facility-scale sources have limited access and time resolution.

Direct high harmonic generation (HHG) driven by intense femtosecond laser pulses can generate multiple harmonic orders throughout the VUV, extreme UV (EUV), and soft X-ray (SXR) spectral regions when very high intensities (>10¹³ W/cm²) are available. However, scaling to higher repetition rates with lower pulse energy is not as simple as focusing tighter to reach the same peak intensities. In a free-focus geometry, small focal spot sizes correspond to impractically short lengths and high gas pressures [23], while for HHG in a hollow capillary waveguide, the confinement loss scales with the inverse cube of the core diameter [24]. HHG using solid [25,26] and liquid [27] targets has recently been investigated with good success. However, only a few papers [27,28] have demonstrated generation into the VUV, and the path to scaling the flux and efficiency to enable applications is still unclear.

Negative curvature antiresonant hollow-core fibers offer an attractive alternative to a simple hollow capillary waveguide, as they use microstructures in the core region to induce interference
effects that confine light to a small diameter with minimal propagation loss [29–31]. Notably, HHG has been demonstrated in a similar antiresonant photonic crystal fiber at 1 kHz repetition rate [32]. Nevertheless, using megahertz (MHz)-repetition-rate fiber lasers, it is still impractical to achieve peak intensities high enough for efficient conversion using direct HHG, without rapid damage to the core microstructures.

Here we utilize a new highly cascaded harmonic generation (HCHG) process to enable the simultaneous production of 13 UV/VUV spectral lines using driving laser intensities well below the threshold required for HHG [33,34]. By focusing two colors (the fundamental and second harmonic of a 10 W average power Yb:fiber laser) into a xenon-filled negative curvature hollow-core fiber, and tuning the xenon pressure to provide optimal phase matching, both even and odd harmonic orders are generated ranging from the 3rd to the 15th harmonic order. We use a peak intensity of approximately $2 \times 10^{12} \text{W/cm}^2$, which is significantly lower than the $\sim 1 \times 10^{14} \text{W/cm}^2$ typically required for efficient HHG [35]. Indeed, at the intensity used in our experiments, the ponderomotive energy of the electron is only 0.1 eV, and the single-atom, single-color ionization probability [36,37] is less than $10^{-15}$—far below the typical HHG regime. As expected for a perturbative cascaded interaction, the flux at each harmonic decreases with increasing harmonic order, but at $\sim 10^{14} - 10^{15} \text{ph/s}$ it rivals synchrotron flux levels [8] for photon energies up to 10.8 eV. Using a model that only includes the third-order nonlinearity ($\chi^{(3)}$), we confirm our hypothesis regarding the cascaded mechanism for harmonic generation and gain a deeper understanding of the physics of HCHG.

2. EXPERIMENT

We start with a prototype ultrafast ytterbium-doped fiber laser producing 10 µJ, 160 fs pulses at 1035 nm with a repetition rate of 1 MHz. The output is then frequency doubled to 518 nm [beta-barium borate (BBO), Eksma Optics] with about 50% efficiency. The 1035 and 518 nm pulses are overlapped in time using a delay stage and focused into a 30 µm diameter negative-curvature fiber (Glo-Photonics PMC-500/700) filled with $\sim 1000$ Torr (20 psia) of xenon gas [Fig. 1(a)]. With 4 W of 1035 nm light and 3 W of 518 nm light entering the fiber (parallel linear polarization), we generate about 200 mW of third harmonic (345 nm) through a degenerate four-wave-mixing process—here two photons of the second harmonic are combined to produce one photon each of the fundamental and the third harmonic. This third-harmonic light can then combine with the fundamental and second-harmonic light to drive additional four-wave-mixing processes in a cascaded series [Fig. 1(b)]. Such cascaded processes have been studied previously, but past studies [38–40] have used Ti:Sapphire lasers with shorter pulses, order-of-magnitude-higher pulse energies, and larger-diameter capillary waveguides. In contrast, we employ a compact fiber laser and a small diameter negative-curvature waveguide to produce harmonics up to the 15th harmonic.

After generation, the VUV light passes through an in-vacuum monochromator in order to select a single wavelength for transmission to the sample. In these experiments, we employed three different monochromator designs. (1) First, we used a prism monochromator to measure the flux of harmonic orders 3–9. In this setup, the VUV beam was collimated using a concave collimating mirror (R = 500 mm), diffracted through a MgF$_2$ prism (60° apex), reflected by a flat wavelength-tuning mirror, and focused with the beam, placed immediately after the exit of the hollow fiber. This dielectric mirror was designed to transmit most of the 1035, 517, and 345 nm light while reflecting all shorter wavelengths. This mirror greatly reduces the heat load on the more sensitive Al + MgF$_2$-coated optics. This also allowed us to direct the 1035, 517, and 345 nm beams outside the vacuum chamber to monitor for optimal fiber coupling and temporal overlap.

Fig. 1. Highly cascaded harmonic generation (HCHG). (a) In our experiment, a 1035 nm laser is frequency doubled, and both fundamental and second-harmonic beams are focused into a xenon-filled, 30 µm core diameter, hollow-core negative-curvature fiber to drive the HCHG process. The resulting VUV light is then spectrally selected by a grating-based monochromator. (b) The HCHG process is the result of numerous four-wave-mixing steps, each combining three photons to generate a higher-energy photon. We note that the photon combinations for H4 and H9 represent just one possible route to each harmonic.
3. RESULTS AND DISCUSSION

Using highly cascaded harmonic generation, we simultaneously generate bright even and odd harmonic orders 3–15 of the bichromatic 1035 and 517 nm driving laser. The measured radiant flux (radiant power) and the corresponding photon flux of each harmonic are recorded in Table 1 and shown in Fig. 2. High-resolution spectra of the eighth and ninth harmonics [Fig. 2(a), shown in red] reveal bandwidths of 40 meV (full width at half-maximum), which is somewhat more bandwidth than the 8 meV bandwidth of the driving laser, indicating that the HCHG process likely produces shorter pulses, or pulses that can be compressed to be shorter, than the driving pulses. This type of temporal shortening is often seen with nonlinear optical processes such as HHG [41].

The HCHG process demonstrated here is a cascaded four-wave-mixing process. This is distinct from HHG, where each harmonic is generated directly from the driving laser using high-wave-mixing process. This is distinct from HHG, where each harmonic is generated directly from the driving laser using high-wave-mixing process such as HHG [41].

Table 1. Observed and Estimated Source Photon Flux for Each Harmonic

<table>
<thead>
<tr>
<th>Harmonic Order</th>
<th>Photon Energy (eV)</th>
<th>Wavelength (nm)</th>
<th>Power (µW)</th>
<th>Measured Photon Flux (ph/s)</th>
<th>Estimated Source Flux from Harmonic (ph/s)</th>
</tr>
</thead>
<tbody>
<tr>
<td>3</td>
<td>3.6</td>
<td>345</td>
<td>230000</td>
<td>4.0  × 10^{17}</td>
<td>4.7  × 10^{17}</td>
</tr>
<tr>
<td>4</td>
<td>4.8</td>
<td>259</td>
<td>6000</td>
<td>7.8  × 10^{13}</td>
<td>1.6  × 10^{16}</td>
</tr>
<tr>
<td>5</td>
<td>6.0</td>
<td>207</td>
<td>2200</td>
<td>2.3  × 10^{15}</td>
<td>4.6  × 10^{16}</td>
</tr>
<tr>
<td>6</td>
<td>7.2</td>
<td>173</td>
<td>610</td>
<td>5.3  × 10^{14}</td>
<td>1.6  × 10^{15}</td>
</tr>
<tr>
<td>7</td>
<td>8.4</td>
<td>148</td>
<td>120</td>
<td>9.1  × 10^{13}</td>
<td>2.7  × 10^{14}</td>
</tr>
<tr>
<td>8</td>
<td>9.6</td>
<td>129</td>
<td>26</td>
<td>1.7  × 10^{13}</td>
<td>6.8  × 10^{13}</td>
</tr>
<tr>
<td>9</td>
<td>10.8</td>
<td>115</td>
<td>4.6</td>
<td>2.6  × 10^{12}</td>
<td>2.1  × 10^{13}</td>
</tr>
<tr>
<td>10</td>
<td>12.0</td>
<td>104</td>
<td>0.036</td>
<td>1.9  × 10^{10}</td>
<td>3.0  × 10^{12}</td>
</tr>
<tr>
<td>11</td>
<td>13.2</td>
<td>94</td>
<td>0.037</td>
<td>1.7  × 10^{10}</td>
<td>3.0  × 10^{12}</td>
</tr>
<tr>
<td>12</td>
<td>14.4</td>
<td>86</td>
<td>0.0079</td>
<td>3.4  × 10^{9}</td>
<td>9.0  × 10^{11}</td>
</tr>
<tr>
<td>13</td>
<td>15.6</td>
<td>80</td>
<td>0.0036</td>
<td>1.4  × 10^{9}</td>
<td>3.0  × 10^{11}</td>
</tr>
<tr>
<td>14</td>
<td>16.8</td>
<td>74</td>
<td>0.00059</td>
<td>2.2  × 10^{8}</td>
<td>5.0  × 10^{10}</td>
</tr>
<tr>
<td>15</td>
<td>18.0</td>
<td>69</td>
<td>0.00035</td>
<td>1.2  × 10^{8}</td>
<td>3.0  × 10^{10}</td>
</tr>
</tbody>
</table>

*Harmonic orders 3–9 were measured using a prism-based monochromator, while harmonics 9–15 were measured using a grating-based Wadsworth monochromator. The estimated photon flux at the source for harmonics 10–15 should be regarded as an order-of-magnitude estimate and is limited by knowledge of the precise reflectivity of the optics.

Second harmonic and fundamental beams [Fig. 1(b)]. The energy conservation is described by

\[ \omega_3 = \omega_2 + \omega_2 - \omega_1, \]

where \( \omega_n \) is the angular frequency of the \( n \)th harmonic, \( \omega_n = n \omega_1 \), and \( \omega_1 \) is the frequency of the driving laser (\( \hbar \omega_1 = 1.2 \text{ eV} \)). The phase mismatch in a gas-filled hollow waveguide [42] is given by

\[ \Delta k \equiv k_1 + k_3 - 2k_2 = 2\pi N \left( \frac{\delta_3}{\lambda_3} + \frac{\delta_1}{\lambda_1} - \frac{2\delta_2}{\lambda_2} \right) - \mu \frac{1}{4\pi a^2} (\lambda_3 + \lambda_1 - 2\lambda_2), \]

where \( \lambda_n \) is the \( n \)th harmonic wavelength, \( k_n \) is the corresponding wavevector, \( N \) is the number density of the gas, \( \delta_n \) is the gas dispersion (related to the refractive index by \( n - 1 = \mathcal{N} \mathcal{S} \)) for the \( n \)th harmonic, \( \mu \) is a mode-dependent constant (2.405 for the lowest-order mode used here [42]), and \( a \) is the diameter of the waveguide. The xenon refractive index has been measured [43] for wavelengths longer than 140 nm. The first term on the right side of Eq. (2) is the pressure-dependent contribution from the gas, and the second term is the pressure-independent term from the waveguide confinement. Typically, there is some pressure that will eliminate the phase mismatch, allowing the third-harmonic generation to be phase matched.

Once the third harmonic is produced, this wavelength can lead to the generation of higher harmonics by the same process. Beyond the fourth harmonic, multiple pathways can lead to the production of each harmonic, for example:

\[ \omega_4 = \omega_3 + \omega_2 - \omega_1, \]

\[ \omega_4 = \omega_3 + \omega_3 - \omega_2, \]

\[ \omega_5 = \omega_4 + \omega_2 - \omega_1, \]

\[ \omega_5 = \omega_4 + \omega_3 - \omega_1, \]

\[ \omega_6 = \omega_5 + \omega_2 - \omega_1, \]

\[ \omega_6 = \omega_5 + \omega_3 - \omega_1, \]

\[ \omega_7 = \omega_6 + \omega_2 - \omega_1, \]

\[ \omega_7 = \omega_6 + \omega_3 - \omega_1, \]

\[ \omega_7 = \omega_6 + \omega_4 - \omega_1, \]

\[ \cdots \]
Each of these processes has an optimal phase-matching condition, and most of these processes are phase matched for some xenon pressure below that needed to phase match third-harmonic generation, which is the first step in the cascaded process. Because we apply pressure only to the front of the 100 mm fiber (around 1000 Torr, tuned to optimize harmonic production) and allow the gas to flow through the fiber into a vacuum, we have a gradient of pressure along the length of the fiber [Figs. 3(a) and 3(b)], estimated from a finite element calculation employing a pressure-dependent conductance. The range of xenon pressures in our fiber allows many of the relevant phase-matching and quasi-phase-matching conditions to be met at some point along the fiber. However, more complicated situations are possible—past work [38,39] indicates that quasi-phase-matching conditions arising from the periodic buildup and decay of intermediate harmonics could be more important than true phase matching for cascaded processes.

To provide a more complete picture of the HCHG process, we perform numerical calculations using the nonlinear Schrödinger equation (NLSE) and implemented in the PyNLO package [44,45]. In these simulations, the pulses are modeled with a sech^2 temporal profile with a full width at half-maximum of 160 fs and a pulse energy of 2 µJ each. The pressure-dependent nonlinear index of xenon is assumed to be 1.1 x 10^-22 m^2/(W bar) [46,47]. The calculations predict output fluxes [Fig. 3(c)] roughly comparable to what we measure (Fig. 2), confirming that our experimental results are consistent with a χ^{(3)}-driven cascaded mixing process. These calculations reveal that the harmonics are produced all along the length of the fiber [Fig. 3(d)], indicating that quasi-phase matching is likely important in the production of bright harmonics. The pressure profile used here was chosen primarily for experimental convenience and may not be ideal for harmonic generation. Further experimentation and more advanced calculations are needed to determine the ideal pressure profile for HCHG.

One advantage of direct HHG is the ability to generate wavelengths extending far beyond the ionization energy of the nonlinear medium where absorption is very high. In contrast, the cascaded harmonic generation shown here could reasonably be expected to cut off at the first harmonic that is strongly absorbed by the medium. However, an additional calculation shown in Fig. 4 reveals that attenuation of harmonics beyond 3ω have only a small effect on the flux of higher harmonics. Moreover, the experimental results shown in Fig. 2 also indicate that absorption of a higher harmonic does not prevent the generation of higher harmonics. The ionization energy of xenon is 12.13 eV, with a high density of Rydberg states in the vicinity of our 10th harmonic at 12.0 eV. As shown in Fig. 2, we observe a lower intensity at this wavelength, yet harmonics up to the 15th are observed. We therefore conclude that the HCHG process has some ability to generate harmonics beyond absorption bands in the medium. Furthermore, in HCHG, we can effectively drive the nonlinear optical process with relatively long-duration pulses, giving us the opportunity to generate harmonics with spectral bandwidth significantly narrower than is typical for HHG.

The HCHG process is a powerful tool for bringing high-flux VUV light to tabletop experiments. To our knowledge, this is the smallest, most energy-efficient source of femtosecond pulses of vacuum-ultraviolet light in the 7–11 eV spectral range. The total
energy consumption of all electronics and power supplies used for this source is approximately 1 kW (an additional ~1 kW is used by water chillers). Nevertheless, our flux for our harmonics in the 6–10 eV spectral range meets or exceeds that of the Advanced Light Source synchrotron (2 × 10^{13} photons/s for similar bandwidth) [48]. Thus, bright multispectral VUV light can be implemented in individual labs, and synchrotron facilities can then be used more efficiently for experiments requiring fully tunable VUV light or high fluxes at higher photon energies.

The VUV light produced by this source is linearly polarized, which is a valuable feature that broadens the scope of possible applications [49]. Any linear polarization for the VUV can be produced, as the polarization of the driving lasers is conserved. Additionally, the use of circularly polarized driving lasers may allow circularly polarized light to be generated directly. Circularly polarized harmonics have been produced from both HHG (counter-rotating fields) [50] and four-wave mixing (seed co-rotating with a single pump) [51]. Since HCHG is simply a cascade of multiple four-wave-mixing steps, it should also produce circularly polarized light when co-rotating fields are used. Future studies will be required to experimentally verify this capability.

This source could enable major advancements in important scientific applications such as time-, spin-, and/or angle-resolved photoemission spectroscopy [1–5,52–54] as well as PIMS [6–8,12]. For example, the high repetition rate of this laser is well suited for ARPES experiments that can suffer from a loss of energy resolution due to space charge effects. Moreover, the tunability of this VUV source is currently being used for PIMS experiments, where it has demonstrated the ability to differentiate between different molecules that have the same mass [55].

4. CONCLUSION

We have demonstrated the generation of high-flux vacuum-ultraviolet light using highly cascaded harmonic generation. Numerical simulations show that the harmonic generation is driven by cascaded four-wave-mixing processes. The observed spectral bandwidths of ~40 meV are ideal for many scientific applications, and narrower spectra down to 1 meV can be obtained using moderately sized monochromators. The process of highly cascaded harmonic generation can bring bright, high-repetition-rate, multispectral VUV light, previously available only at large facilities, to laboratory tabletops.

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Disclosures. DDH, SJB, DGW, MSK, SRD, JJR, HCK: KMLabs (E,P). MMM, HCK, SJB: KMLabs (I). KMLabs uses this HCHG technology to build VUV laser systems.

REFERENCES


