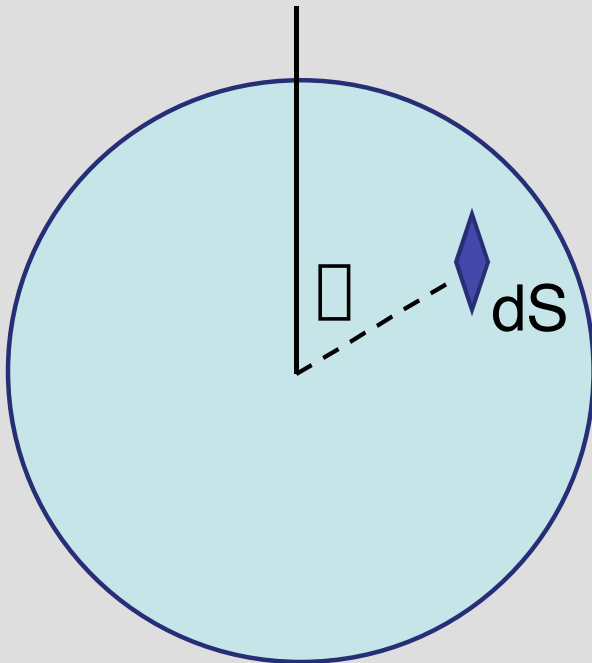


# The equation of radiative transfer

How does the intensity of radiation change in the presence of emission and / or absorption?

## Definition of solid angle and steradian



Sphere radius  $r$  - area of a patch  $dS$  on the surface is:

$$dS = r d\theta \cdot r \sin\theta d\phi \equiv r^2 d\Omega$$

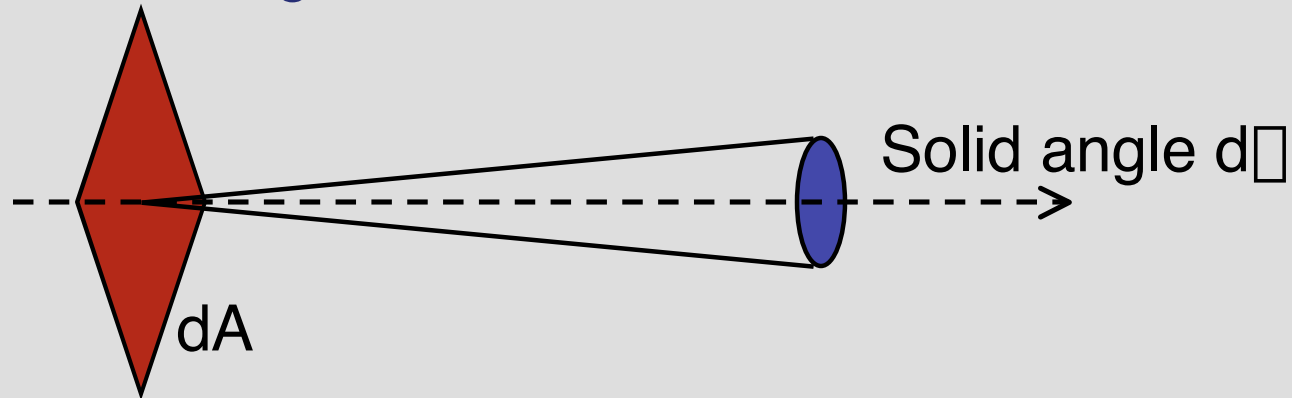
$d\Omega$  is the solid angle subtended by the area  $dS$  at the center of the sphere.

Unit of solid angle is the steradian.

$4\pi$  steradians cover whole sphere.

## Definition of the specific intensity

Construct an area  $dA$  normal to a light ray, and consider all the rays that pass through  $dA$  whose directions lie within a small solid angle  $d\Omega$ .



The amount of energy passing through  $dA$  and into  $d\Omega$  in time  $dt$  in frequency range  $d\nu$  is:

$$dE = I_\nu dA dt d\Omega d\nu$$



**Specific intensity** of the radiation.

Compare with definition of the flux: specific intensity is very similar except it depends upon direction and frequency as well as location.

Units of specific intensity are:  $\text{erg s}^{-1} \text{ cm}^{-2} \text{ Hz}^{-1} \text{ steradian}^{-1}$

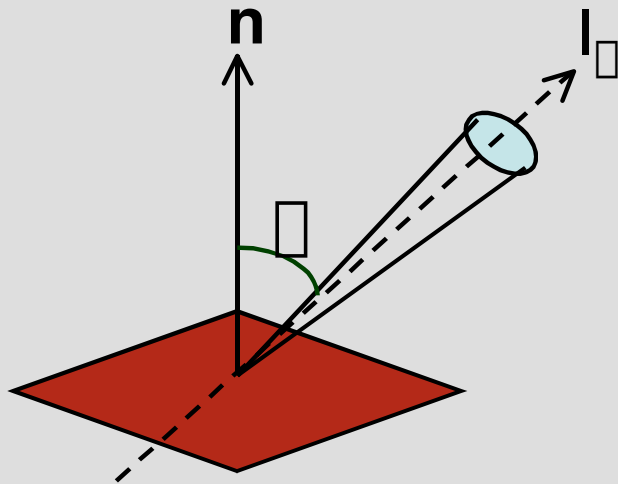


Same as  $F_{\lambda}$

Another, more intuitive name for the specific intensity is **brightness**.

Simple relation between the flux and the specific intensity:

Consider a small area  $dA$ , with light rays passing through it at all angles to the normal to the surface  $\mathbf{n}$ :



If  $\theta = 90^\circ$ , then light rays in that direction contribute **zero** net flux through area  $dA$ .

For rays at angle  $\theta$ , foreshortening reduces the effective area by a factor of  $\cos(\theta)$ .

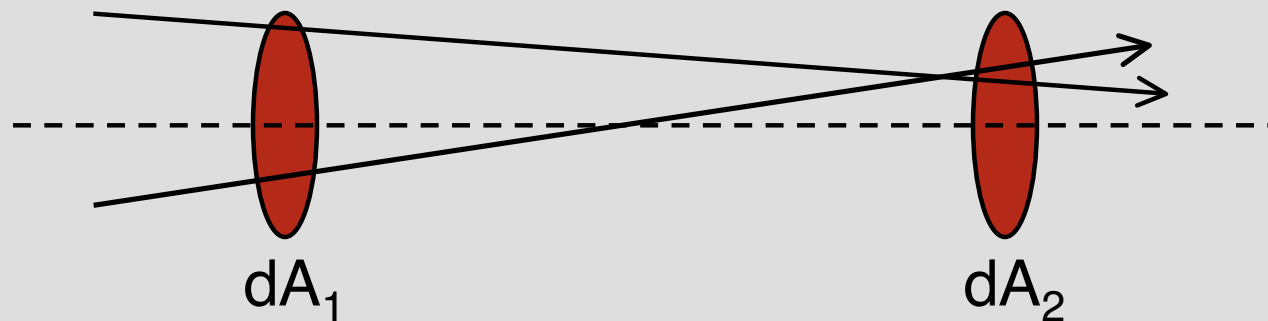
Hence, net flux in the direction of  $\mathbf{n}$  is given by integrating (the specific intensity  $\times \cos \theta$ ) over all solid angles:

$$F_{\mathbf{n}} = \int I_{\theta} \cos \theta d\Omega$$

Note: to actually evaluate this need to express  $d\Omega$  in terms of  $d\theta$  and  $d\phi$  as before.

## How does specific intensity change along a ray

If there is no emission or absorption, specific intensity is just **constant** along the path of a light ray. Consider any two points along a ray, and construct areas  $dA_1$  and  $dA_2$  normal to the ray at those points. How much energy is carried by those rays that pass through **both**  $dA_1$  and  $dA_2$ ?



$$\left. \begin{aligned} dE_1 &= I_{\square_1} dA_1 dt d\square_1 d\square_1 \\ dE_2 &= I_{\square_2} dA_2 dt d\square_2 d\square_2 \end{aligned} \right\} \text{ where } d\square_1 \text{ is the solid angle subtended by } dA_2 \text{ at } dA_1 \text{ etc}$$

The same photons pass through both  $dA_1$  and  $dA_2$ , without change in their frequency. Conservation of energy gives:

$$dE_1 = dE_2 \quad - \text{equal energy}$$

$$d\Omega_1 = d\Omega_2 \quad - \text{same frequency interval}$$

Using definition of solid angle, if  $dA_1$  is separated from  $dA_2$  by distance  $r$ :

$$d\Omega_1 = \frac{dA_2}{r^2}, \quad d\Omega_2 = \frac{dA_1}{r^2}$$

Substitute:

$$I_{\Omega_1} dA_1 dt d\Omega_1 d\Omega_1 = I_{\Omega_2} dA_2 dt d\Omega_2 d\Omega_2 \quad dE_1 = dE_2$$

$$I_{\Omega_1} dA_1 dt d\Omega_1 \frac{dA_2}{r^2} = I_{\Omega_2} dA_2 dt d\Omega_2 \frac{dA_1}{r^2} \quad d\Omega_1 = d\Omega_2$$

$$\underline{I_{\Omega_1} = I_{\Omega_2}}$$

Conclude: specific intensity remains the same as radiation propagates through free space.

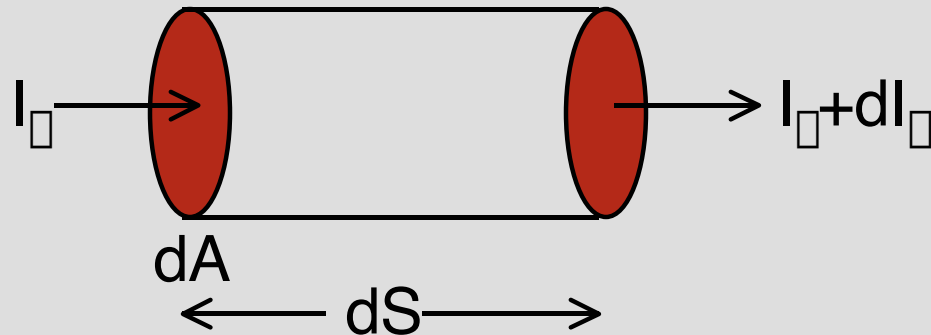
Justifies use of alternative term 'brightness' - e.g. brightness of the disk of a star remains same no matter the distance - **flux** goes down but this is compensated by the light coming from a smaller area.

If we measure the distance along a ray by variable  $s$ , can express result equivalently in differential form:

$$\frac{dI_{\square}}{ds} = 0$$

## Emission

If the radiation travels through a medium which is itself emitting radiation, that will add to the energy:



Spontaneous **emission coefficient** is the amount of energy emitted per unit time, per unit solid angle, per unit frequency interval, and per unit volume:

$$dE = j_\nu dV d\Omega dt d\nu$$

In going a distance  $ds$ , beam of cross-section  $dA$  travels through a volume  $dV = dA \times ds$ .



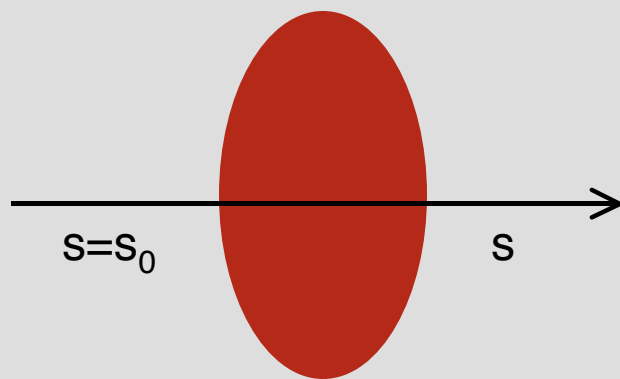
Change (increase) in specific intensity is therefore:

$$dI_{\lambda} = j_{\lambda} ds$$

Equation of radiative transfer for pure emission becomes:

$$\frac{dI_{\lambda}}{ds} = j_{\lambda}$$

If we know what  $j_{\lambda}$  is, can integrate this equation to find the change in specific intensity as radiation propagates through the gas:



$$I_{\lambda}(s) = I_{\lambda}(s_0) + \int_{s_0}^s j_{\lambda}(s') ds'$$

i.e. add up the contributions to the emission all along the path.